

# BEAM-BEAM INDUCED ORBIT EFFECTS AT LHC

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## Abstract

For high bunch intensities the beam-beam force is strong enough to expect orbit effects if the two beams do not collide head-on but with a crossing angle or with a given offset. As a consequence the closed orbit changes. The closed orbit of an unperturbed machine with respect to a machine where the beam-beam force becomes more and more important has been studied and the results are presented in this paper.

## INTRODUCTION

In LHC (Large Hadron Collider) the beam-beam electromagnetic force is experienced as a localized, periodic distortion when the two beams interact with each other at the collision points. This force is most important for high brightness beams which is the case at LHC, and can be classified in two types: head-on and long-range. The effects of the beam-beam force manifest themselves in very different ways. In this paper the closed orbit effects due to long-range beam-beam interactions are studied. Long-range interactions distort the beams much less than head-on interactions. However, there is a large number of them due to the large number of bunches per beam (2808 bunches per beam in nominal LHC). In nominal conditions, up to 30 long-range interactions per experiment have to be expected. For the first time, experimental data is presented in order to validate the simulation studies about the beam-beam effects at the LHC [1], [2], [3] performed during the past years.

## LONG-RANGE BEAM-BEAM KICK

When the beams are separated, as it is the case for long-range beam-beam interactions, the beams will exert a kick to each other whose coherent dipole component leads to orbit changes. The change in angle (kick) can be computed with the following equation [1]

$$\Delta r' = -\frac{2Nr_0}{\gamma} \cdot \frac{1}{r} \cdot \left[ 1 - \exp\left(-\frac{r^2}{4\sigma^2}\right) \right] \quad (1)$$

where  $r$  is the beam separation,  $\sigma$  the beam size at the interaction point and  $\gamma$  is the relativistic Lorentz factor.

## Analysis and Results

During a dedicated machine study period with only one bunch circulating per beam, a horizontal orbit scan in steps of  $100\ \mu\text{m}$  was performed in ATLAS (IP1) and a vertical one in CMS (IP5) to measure the orbit kick due to the beam-beam force. The scan went far enough to reach the non-linear regime of the force. To avoid additional effects

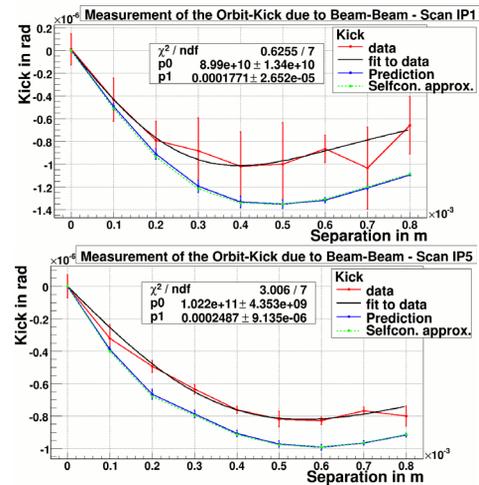


Figure 1: Measurement and prediction of the orbit kick due to beam-beam force during an orbit scan. The vertical axis is given in units of  $10^{-6}$  rad and the horizontal axis in units of  $10^{-3}$  m.

arising from the trim of the separation bumps, only beam 1 was moved and beam 2 was untouched. In this way beam 2 is only perturbed by the beam-beam kick and a clean signal should be measured.

The result is illustrated in Fig. 1, where the red line indicates the kick determined by an orbit correction of the taken data and the solid blue line was calculated using Eq. 1.

For the calculation, the values of  $\sigma$  and  $N$  were varied at each separation step according to the measured values. The shape of the measured curve fits to the expectation, but the measured strength is lower than the prediction. The black line is a fit to the data using Eq. 1, the intensity  $N = p_0$  and the beam size  $\sigma = p_1$  were left as free but constant parameters. The obtained value for  $\sigma$  is in agreement with the measurement, but the obtained  $N$  is too small compared to the measured one, thus the origin of the discrepancy must be different. Therefore, the green dashed line tries to approximate the self-consistent effect which is not considered in Eq. 1 and arises from the effect of the interaction itself. If the orbit is changed due to a bunch crossing, the position of the beams will be different after one turn, which implies a variation of the force, leading again to a different orbit until an equilibrium is found. The dashed line shows that this effect is too small to explain the gap. Further investigation is needed to identify the discrepancy.

Fig. 2 shows the beam position at IP1 and IP5 during the scans as a function of the separation in red, where the error bars show the standard deviation. A prediction was calculated and simulated with the MADX program. Also here a discrepancy between prediction and measurement is found.

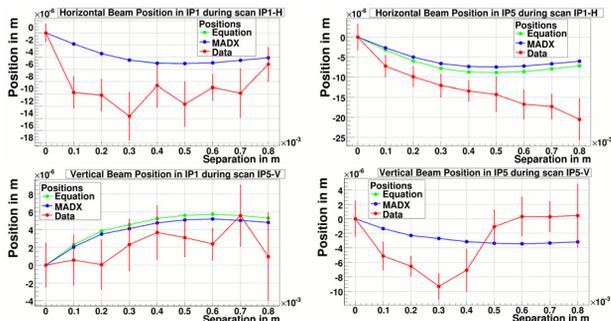


Figure 2: Measurement and prediction of the position at IP1 and IP5 due to beam-beam force during an orbit scan. The vertical axis is given in units of  $10^{-6}$ m and the horizontal one in units of  $10^{-3}$ m.

Nevertheless, the position at the IPs were interpolated from the data of the two closest Beam Position Monitors (BPM) on the left and right hand side of the IP only. Those BPMs have a resolution of around  $5 \mu\text{m}$  which is already larger than the standard deviation and the observed discrepancy. As expected, the shape of the curves agree with the prediction (except for the beam position in IP5 during the scan of IP5, which needs further investigation). However, the source of the gap will be explained when the gap in Fig.1 is understood.

### BUNCH-BY-BUNCH ORBIT DIFFERENCES

The study of the orbit effects due to the coherent dipolar kick of the beam-beam force performed in the previous section shows that the expected orbit variations are in the order of a few micrometers. It is interesting to analyse this effect on a bunch-by-bunch basis, since the so-called PACMAN-effect introduces differences bunch-to-bunch: bunches at the head or tail of a bunch train encounter fewer long-range interactions since they cross empty buckets. This leads to differences in the long-range orbit kicks for those bunches w.r.t to the core of the train [3]. Three ways of measuring the position variations are analysed in the following.

#### Beam Position Monitors

An exhaustive analysis of the data measured by the LHC BPM during a typical luminosity run was performed to assess the feasibility of those instruments to measure closed orbit deviations in bunch-by-bunch mode in the order of a few micrometers.

From the analysis one can conclude that the BPM system has a bunch-by-bunch orbit measurement able to give a relative bunch-to-bunch resolution in the  $5 \mu\text{m}$  range. Yet the non-linearity in the bunch-by-bunch mode for different global positions limits this resolution to  $\pm 50 \mu\text{m}$  when comparing along the train for different mean positions. The current LHC BPM system, hence, does not have sufficient linearity or resolution to resolve the bunch-by-bunch orbit variations at the few micrometer level expected from beam-beam interaction orbit effects.

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### ATLAS Luminous Region Reconstruction

The LHC experiments' vertex detectors, however, have the required resolution to study the beam-beam induced orbit differences between bunches. The analysis presented in the following uses the microvertex detector of the ATLAS experiment. A total of five luminosity fills with different filling schemes have been analysed to investigate the dependence on the number of bunches.

In Fig. 3 the reconstructed luminous region in the vertical

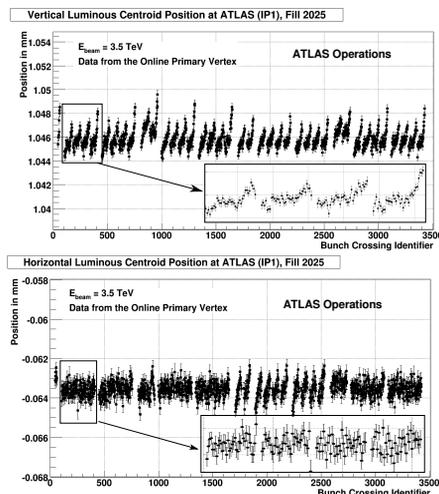


Figure 3: Luminous centroid position in ATLAS. Top (bottom) Figure shows the vertical (horizontal) luminous region position and a zoom over the first four trains of 36 bunches.

(crossing-angle) plane and horizontal (separation) plane of ATLAS is plotted as a function of the bunch crossing identifier for fill 2025 with 1380 bunches spaced by 50 ns. As predicted by the simulations [3], the orbits at the interaction point of all bunches are slightly different and it is more evident for PACMAN bunches in the vertical plane.

The luminous centroid reconstruction can only give the convoluted position of both beams. But since the orbits of both beams are almost identical at the interaction point in the vertical plane (simulations [3] show that they are slightly different due to the intensity variations bunch-by-bunch, of the order of 15% in this fill), most of the bunches collide head-on, although not all of them in the central orbit. In particular, the PACMAN bunches show offsets w.r.t. the core of the train of  $(2.0 \pm 0.3) \mu\text{m}$  due to the different number of long-range interactions. There are two types of them, the ones close to a 8 bunches gap (between trains), and the ones close to a 36/56 bunches gap (between groups of 2 or 4 trains). In Figure 3 a zoom over the first  $4 \times 36$ -bunch-trains is illustrated. The bunches at the core of the train show a different structure than the bunches close to a gap. This is due to the different number of long-range interactions. While the bunches at the core of the train have on average of 74 long-range interactions, the PACMAN bunches have in the order of 10 long-range less. The different number of long-range interactions implies a different amount of coherent dipolar kicks, and therefore dif-

ferent orbits. The effect is even more clear for the PACMAN bunches close to the 36/56 bunches gap, those have on average 40 long-range interactions less and the orbit difference is, consequently, bigger.

The orbit offset due to beam-beam kicks in the vertical plane is  $(5.0 \pm 0.3) \mu\text{m}$  peak to peak. This is a factor 2.5 bigger than the offsets presented in the simulations of Ref. [3]. Nevertheless, it has to be taken into account that the simulations are done for nominal LHC parameters, i.e. 2808 bunches spaced by 25 ns, therefore the bunch pattern is different. As a continuation of this study, new simulations will be done with the parameters of this fill to be able to assess the data in a more quantitative way.

The orbit offsets in the horizontal plane of IP1 do not show a particular structure, as in the case of the vertical plane. If the orbit could be reconstructed for both beams separately, simulations predict that the position of the beams are symmetric to the central orbit and offset by a few micrometers (see Ref. [3] Fig. 14 to 16). The PACMAN bunch effect appears as well, but in the horizontal plane it arises from the long-range interactions in the vertical (crossing) plane of IP5. Since the effect in this plane is global and travels through the ring it appears with a different shape as the local effect in the crossing angle plane of the considered IP. But since we are using the luminous region, only the convolution of both beams is visible, and the effect is cancelled because of its symmetry. There are, however, some trains which show a structure compatible with the beam-beam effect of the vertical plane. The source of this structure should come from extra coherent dipole kicks happening in the horizontal plane somewhere else for those bunches, but this effect needs further investigation.

No dependence on the number of bunches was found for the filling schemes analysed with the same number of long-range interactions, as could be expected.

### Luminosity Optimization

During every luminosity run, once the beams are brought into collisions, every interaction point undergoes a luminosity scan in the horizontal and vertical plane to find the full beam overlap and therefore the maximum luminosity. This process is done integrating the luminosity over the whole beam, and assuming an average position over all bunches. However, since we have demonstrated that there are bunch-by-bunch orbit differences, it is interesting to look at the variation in the maximum of the luminosity as a function of the bunch number. Figure 4 displays the relative position of the two colliding bunches where the maximum luminosity is found during the scan. The plots show IP1 in the horizontal plane (top plot) and vertical plane (bottom plot).

To understand this structure, the position of each beam should be reconstructed. Since this is not possible with the BPMs, simulations have to be referenced again [3] Fig. 14-16: the PACMAN bunches position is offset w.r.t. the core of the train and have a mirror-like structure to the opposite beam. When the scan is performed, the luminosity

maximum is first found for those bunches and afterwards for the core of the train. Thus, in Fig. 4 (top) the gaussian fit mean position is closer to the zero displacement for the PACMAN bunches.

No structure is visible in the vertical plane, because in IP5 the vertical plane is the separation plane where there are no long-range effects. Moreover, the local long-range effects in IP1 can not be seen as in Figure 3, since as explained above the orbits for both beams are nearly identical and as both beams are moved symmetrically in opposite directions the effect is not visible using this measurement method, since the separation is the same for all bunches.

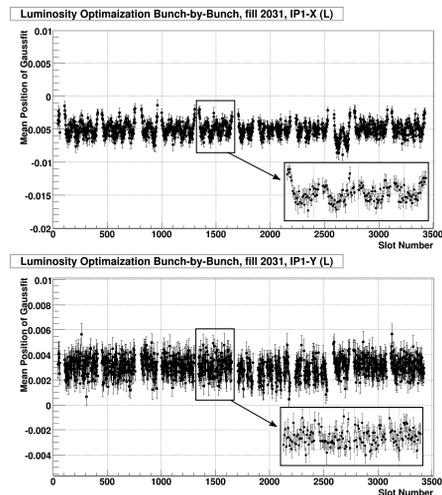


Figure 4: Position of the maximal luminosity (total bunch overlap) as a function of the bunch number. Top (bottom) Figure shows the horizontal (vertical) mean position of the luminosity maximum in units of mm w.r.t. the beam position before the scan.

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### REFERENCES

- [1] Tatiana Pieloni, "A Study of Beam-Beam Effects in Hadron Colliders with a Large Number of Bunches", These no 4211 (2008).
- [2] Werner Herr, "Beam-Beam Interactions", CAS - CERN Accelerator School: Intermediate Course on Accelerator Physics, Zeuthen, Germany, September 2003, CERN-2006-002, p. 379ff.
- [3] Werner Herr, "Features and implications of different LHC crossing schemes", LHC Project Report 628, Geneva, 2003.