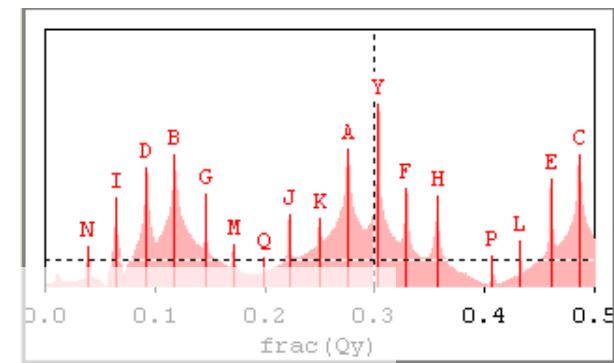
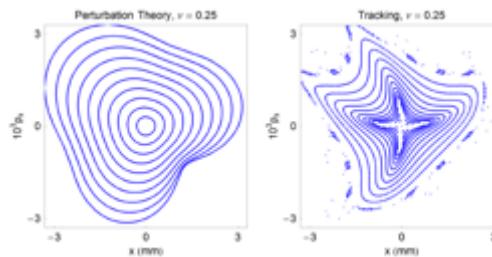
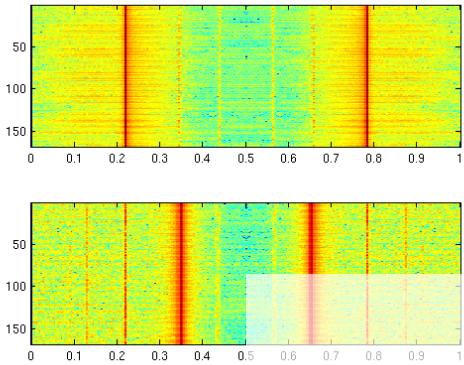
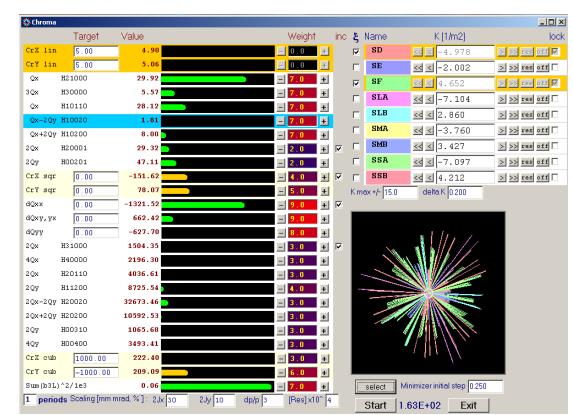
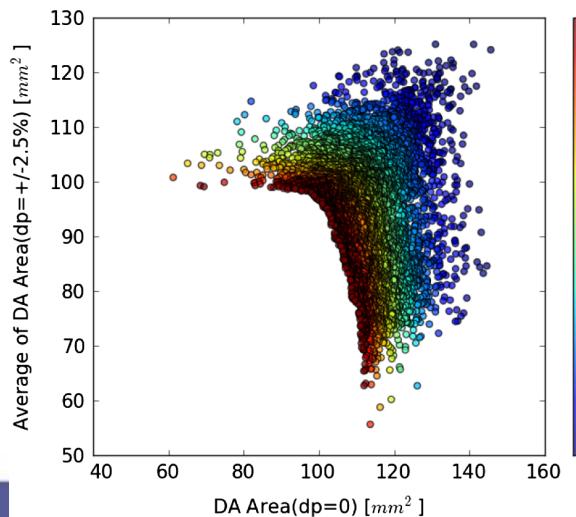
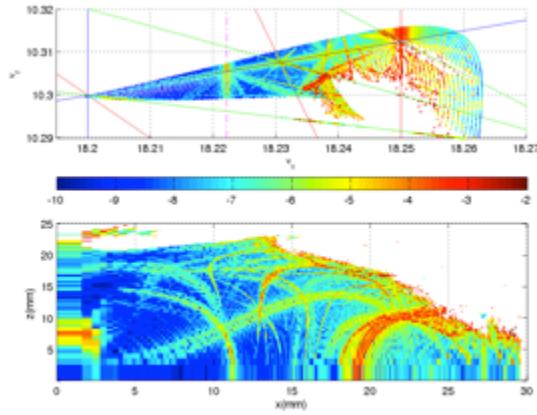


Methods and Tools to Simulate and Analyze non-linear beam Dynamics in Electron Storage Rings



Laurent S. Nadolski

Synchrotron SOLEIL, Gif-sur-Yvette, France



Layout

- **Challenging requirements for today and tomorrow synchrotron light sources**
- **Tools**
 - Tracking codes for storage rings
 - Frequency Map Analysis (FMA) like analysis
 - Momentum aperture & Touschek life time
- **Optimization methods**
 - New optimization method via genetics algorithms
 - Resonance Driving Terms (RDTs)
 - Tuning multipole magnets: sextupole and higher order
- **Perspectives**

Requirements for light sources

- **Low emittance storage ring**
 - Route to diffraction limit in H-plane
 - Complex lattices (TME, MBA lattices)
 - See Dr. Shimosaki's talks
 - Damping wiggler

Storage ring	Energy (GeV)	H-emittance
Petra III	6 Gev	1 nm.rad
PEPX	4.5 GeV	0.1 nm.rad
NSLS II	3 GeV	0.6 nm.rad
MAX IV	3 GeV	< 0.3 nm.rad
Spring 8 upgrade	6 GeV	10 pm.rad
APS upgrade	7 GeV	15 pm.rad

- A **Dynamic Accelerator**:
- **High ratio of straight sections to hosted tens of insertion devices**
 - Apple II, III, Figure 8
 - SC, CRYO undulator/wiggler
 - Controlled polarization undulator
 - Fast switching of B-field (ms scale)

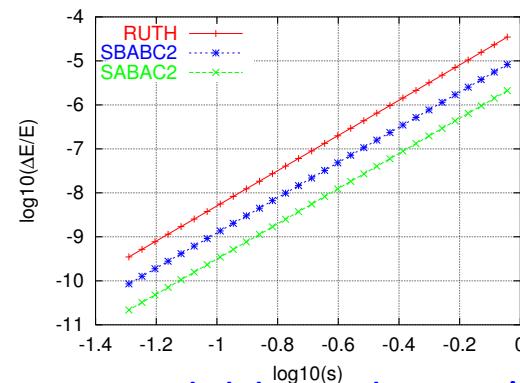
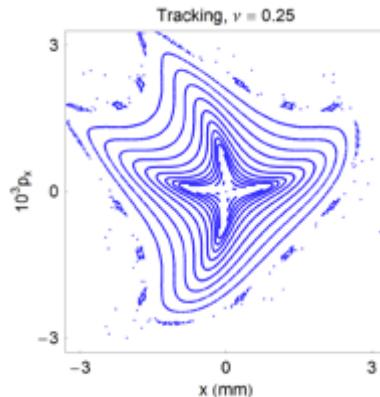
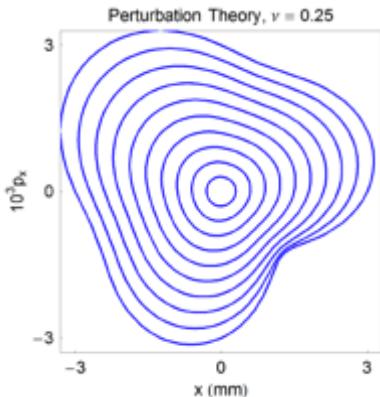
- **High brilliance, flux**
 - Diffracted limited in both planes
- **Taylored filling pattern**
 - Short vs long bunches
 - Low alpha lattice
 - Multi-beam facility (MLS, SOLEIL)
 - M. Ries' s talk today
- **Large on dynamics aperture (15-20mm)**
 - Waiting for on-axis injection
- **Large momentum aperture (2-6%)**
 - Touschek lifetime
 - Reduce beam losses
 - Reduce activation
 - Reduce running cost
- **Top-up operation (MTBF 1 week)**
- **Ultra-low coupling operation**

Figures of merit

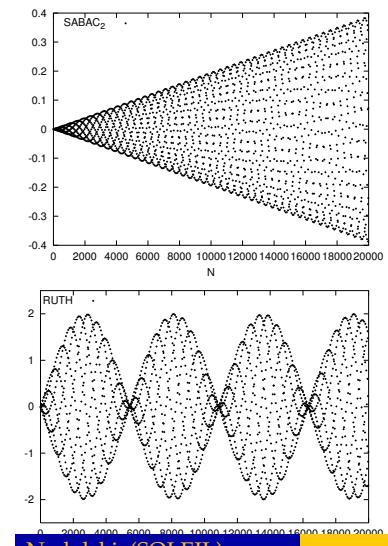
- **How to get there?**
 - High horizontal phase advance
 - Strong focusing lattices, low dispersion function
 - Strong chromatic sextupoles
 - Strong chromatic aberrations
 - Many sextupoles
 - Families
 - Individually powered
 - Higher multipoles
 - **Octupole** (Max IV)
 - Decapole (soon or later)
- **20 years after the first 3GLSs**
 - shift from simple lattice (ALS: 2 sextupole families, TBA based lattice) to highly complex and multi-parameter lattices
- **Optimization of non-linear parameters**
 - Amplitude/momentum tune dependence
 - Non-linear Twiss functions
 - Robust to IDs freely controlled by users

Tracking codes

- Long term tracking based on symplectic integrators
 - Implicit or explicit schemes
- Popularized by Ruth and Forest 1983-1990, use of Lie Algebra (Neri, 1988), Yoshida techniques (1990), Channel and Scovel (1990), McLachlan (1995), Sanz-Serna (1998), Laskar integrators (2001)
 - Preserves energy, bounded errors,
 - Phase stability
 - Used by MADX/PTC, Tracy, OPA, LEGO, ELEGANT, etc.



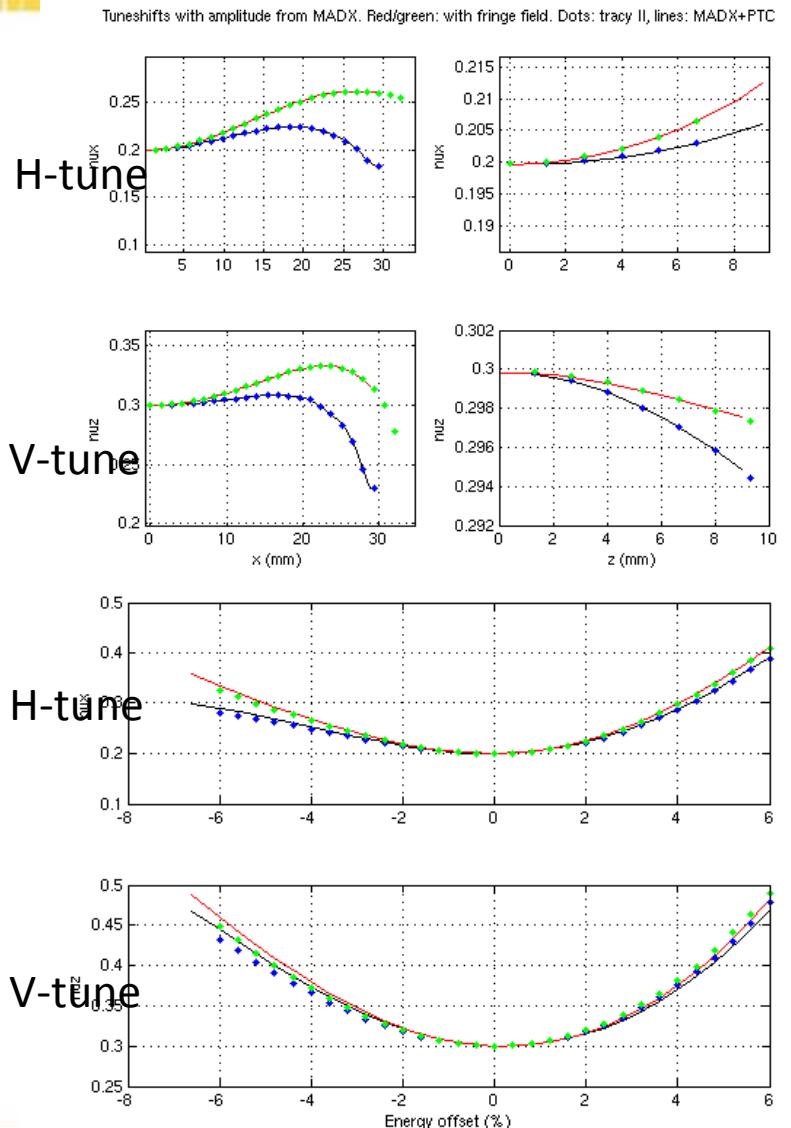
L.S. Nadolski et al. EPAC'02



What is included in the Model?

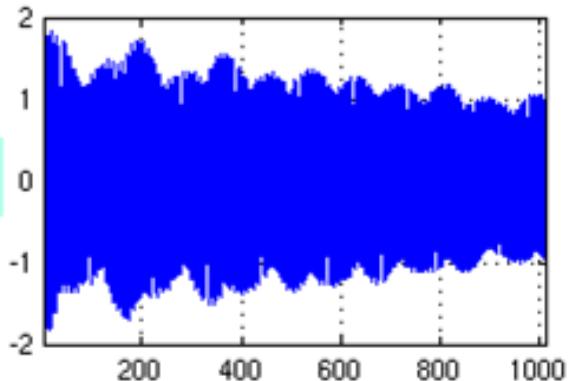
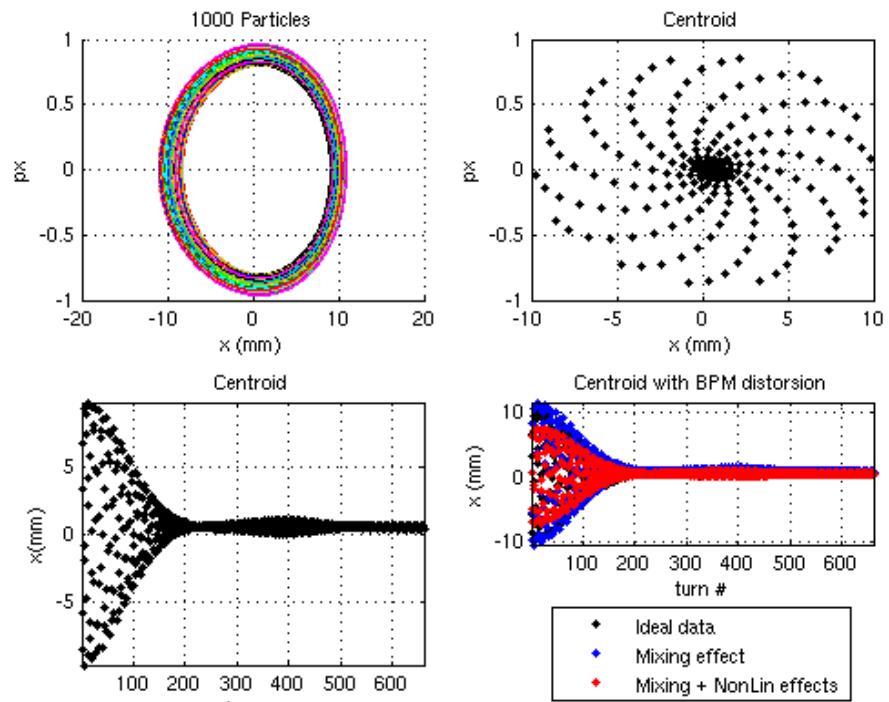
Tracy II & MADX/PTC track exact (SOLEIL case)

- Systematic multipole errors
 - Large momentum acceptance, large dispersion function → high order multipoles
- From magnetic measurements:
 - Add true m-poles (both systematic and non systematic)
 - Dipole: fringe field, gradient error, edge tilt errors
 - Quad.: fringe field
- Beam based
 - Multipoles deduced from turn by turn measurement, off-axis field integrals
 - Coupling errors
- Insertion devices
 - Taylor expansion
 - Radia kick maps
 - Sorting magnets: Genetics algo.
- Collective effects



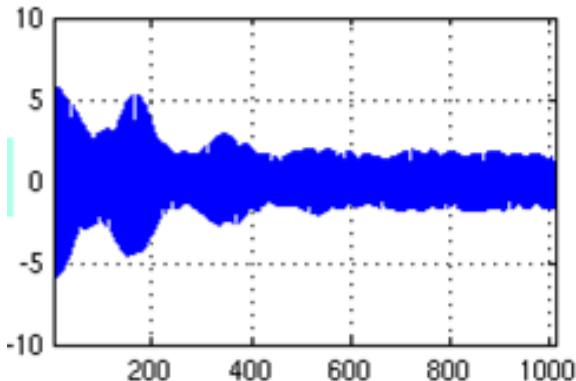
Turn by turn (TbT)data beam smearing

- TbT BPM precision:** $10\mu\text{m}$ $\sim 10 \text{ mA}$: limitative factor Frequency shift
- Algo. to precisely determine tune loose their precision *R. Bartolini et al. Part. Acc. 55, 247 (1995)*
- Lines excited by resonance of order $(m+1)$ decohere m times faster than the tune line. *R. Thomas, PHd Thesis (2003)*
- Gain, coupling correction (LOCO based)



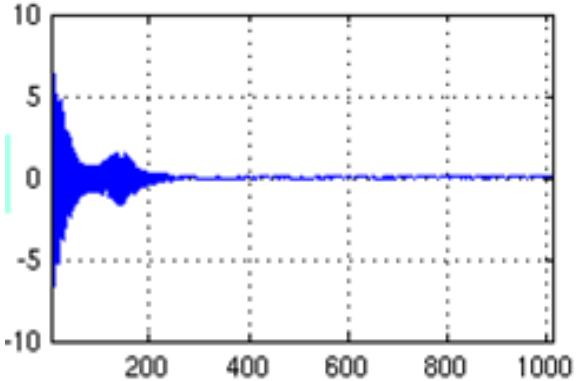
V (kicker H) : 2100 V

x (BPM) : 3.6 mm



9000 V

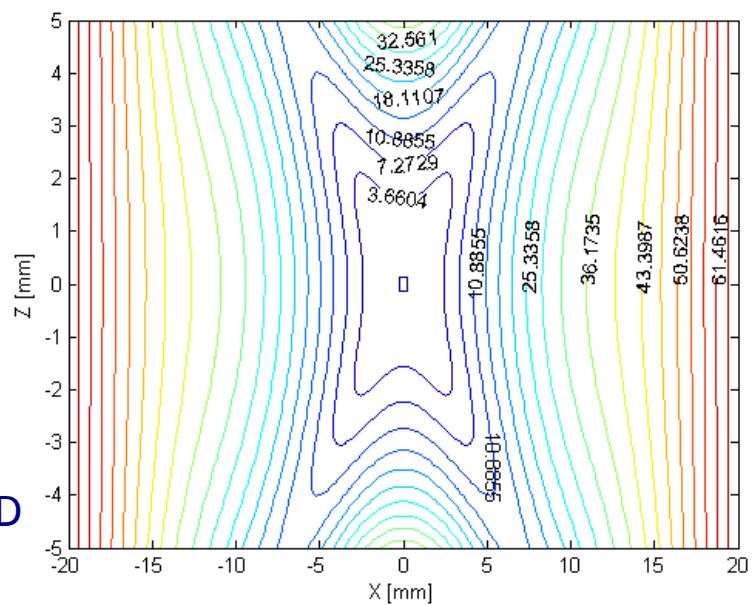
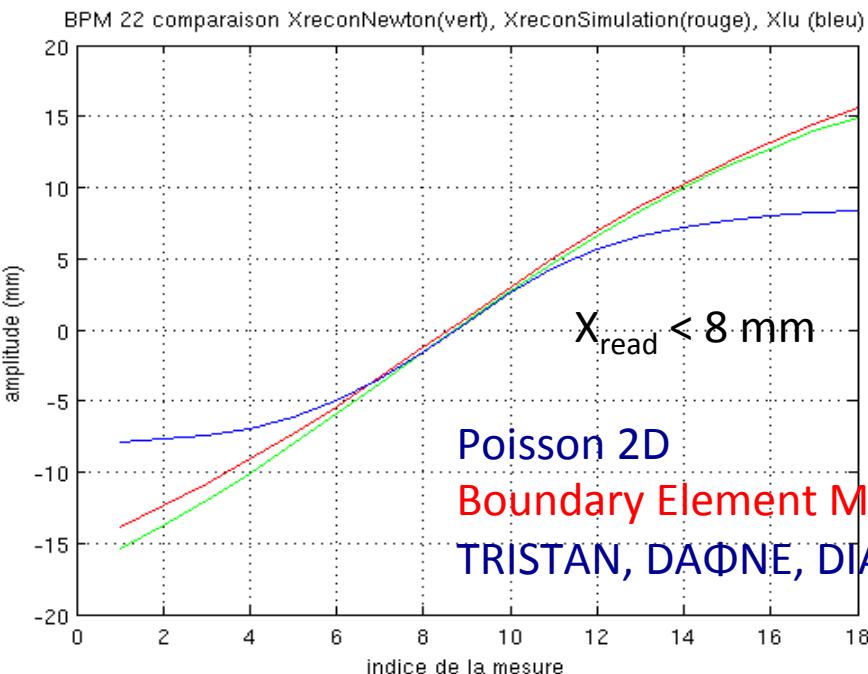
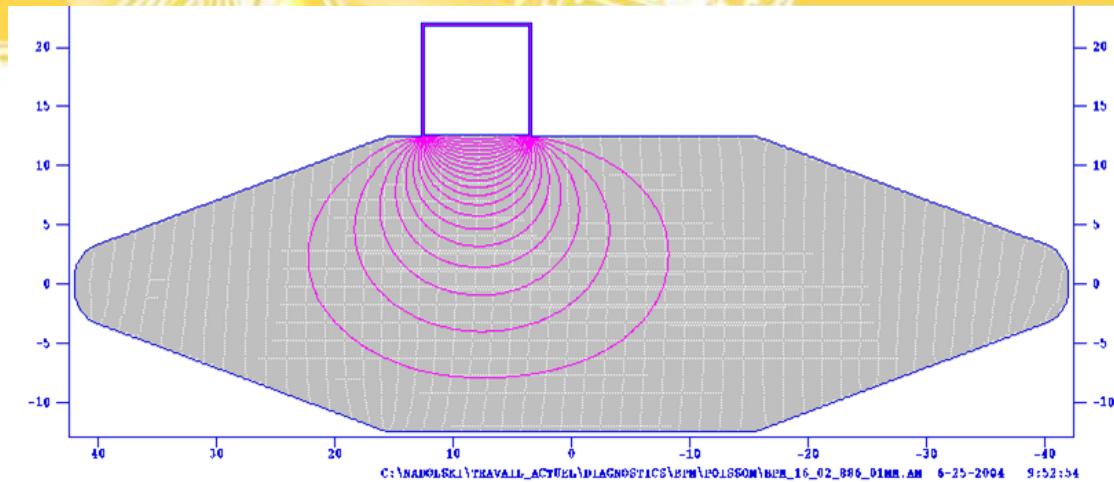
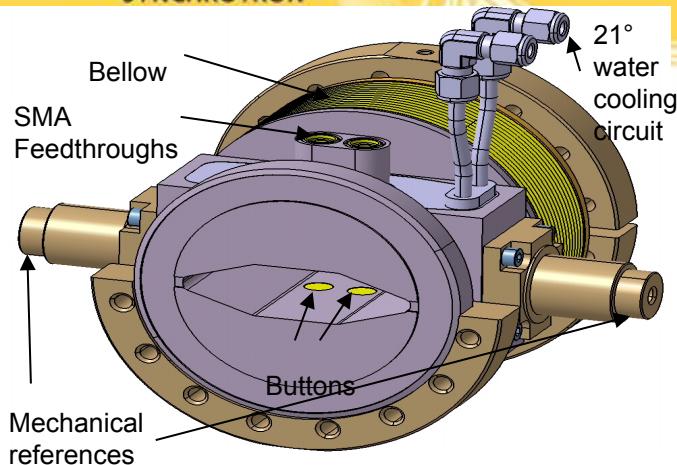
15.2 mm



12100 V

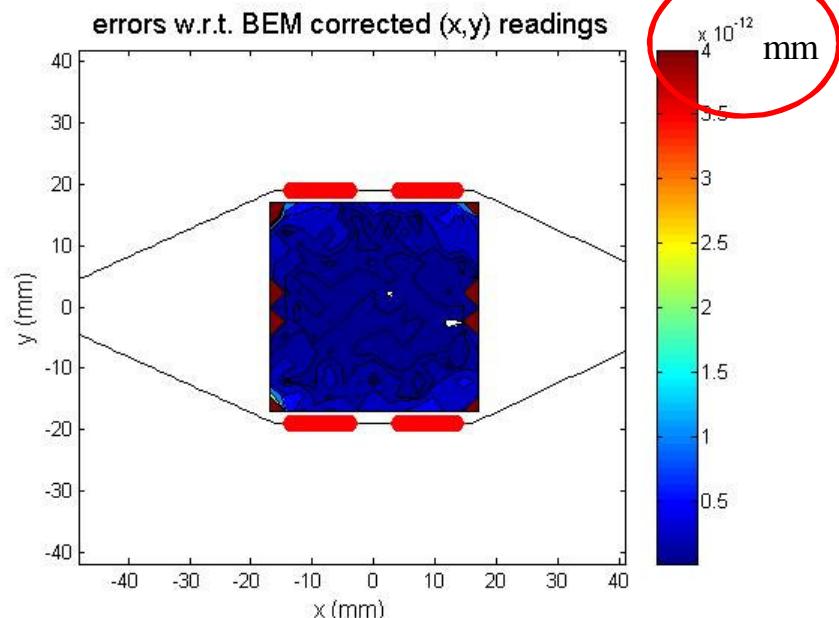
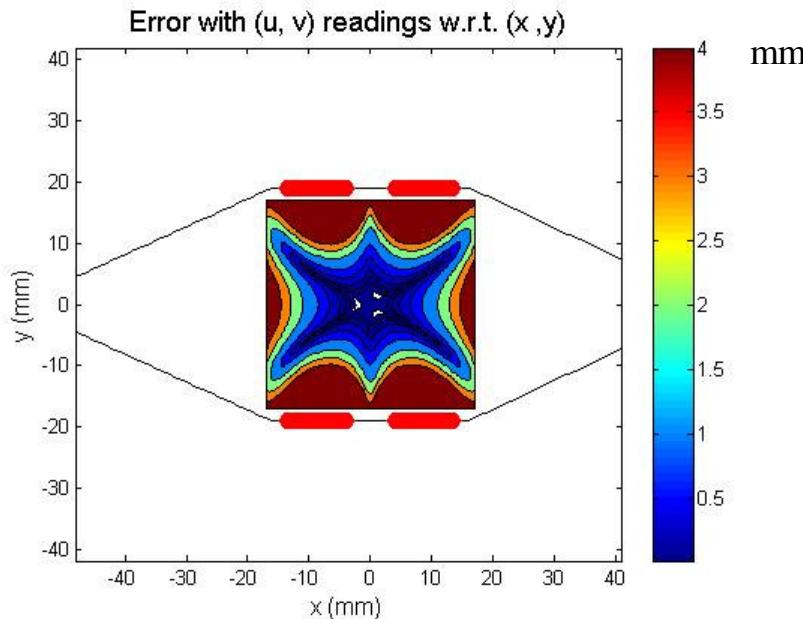
20.6 mm

BPM nonlinear response: SOLEIL case



BPM signal reconstruction

R. Bartolini and J. Rowland, DLS internal note, AP-SR-REP-0171 (March 2010)



Numerical inversion of the map $(x, y) \rightarrow (A, B, C, D)$ given by the electrostatic model

Use (x, y) as fit parameter to reproduce the four buttons readings (A, B, C, D)

3 buttons are sufficient to converge to machine precision

Courtesy of R. Bartolini

Frequency Map Analysis

Laskar A&A1988, Icarus1990 NATO-ASI 1996

- Construction of frequency map

$$\mathcal{F}_\tau : \begin{matrix} \mathbb{R}^n & \longrightarrow & \mathbb{R}^n \\ q|_{p=p_0} & \longrightarrow & \nu \end{matrix}$$

with high precision: $\frac{1}{\tau^4}$ for Hanning Filter

- Refined Fourier technique
- Fast convergence (reduce tracking time)
- Give a global view of the transverse dynamics in a 2D map
- Mapping between DA/tune space using diffusion index

- Determination of tune diffusion vector

$D|_{t=\tau} = \nu|_{t \in (0, \tau/2]} - \nu|_{t \in (\tau/2, \tau]}$
and construction of diffusion map

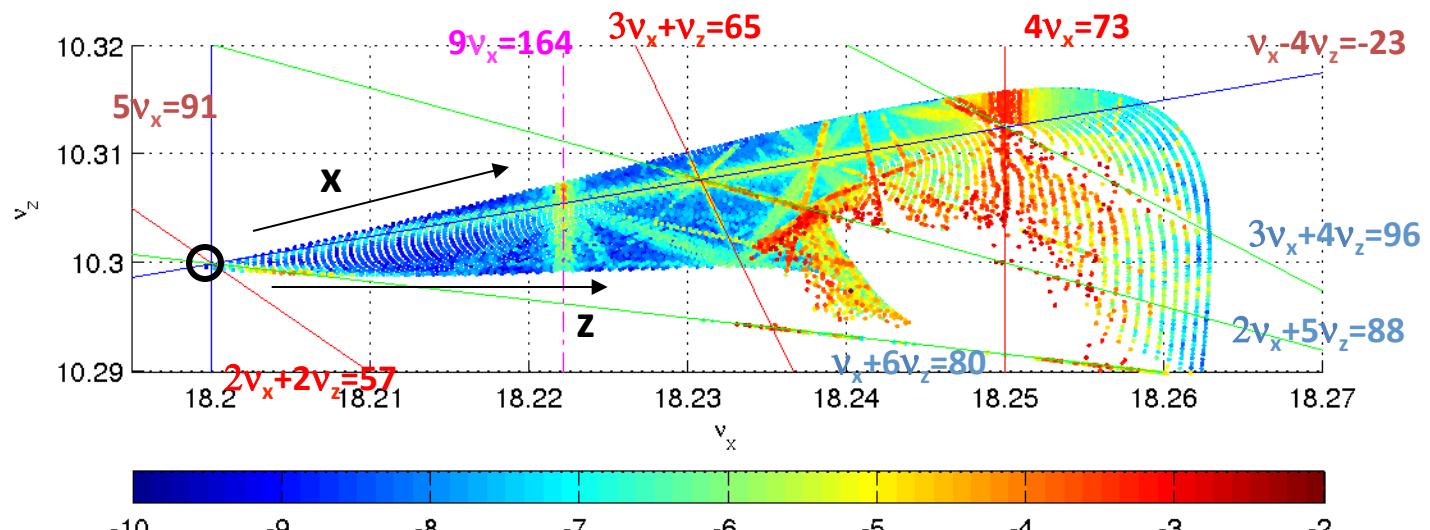
$$\mathcal{D}_\tau : \begin{matrix} \mathbb{R}^n & \longrightarrow & \mathbb{R}^n \\ q|_{p=p_0} & \longrightarrow & D \end{matrix}$$

Does not provide a way to optimize

- Determination of resonance driving terms associated with amplitudes $a_{j,k}$
- Bengtsson PhD thesis CERN88-05

On-momentum Dynamics --Working point SOLEIL lattice (18.2,10.3)

Bare lattice
(no errors)



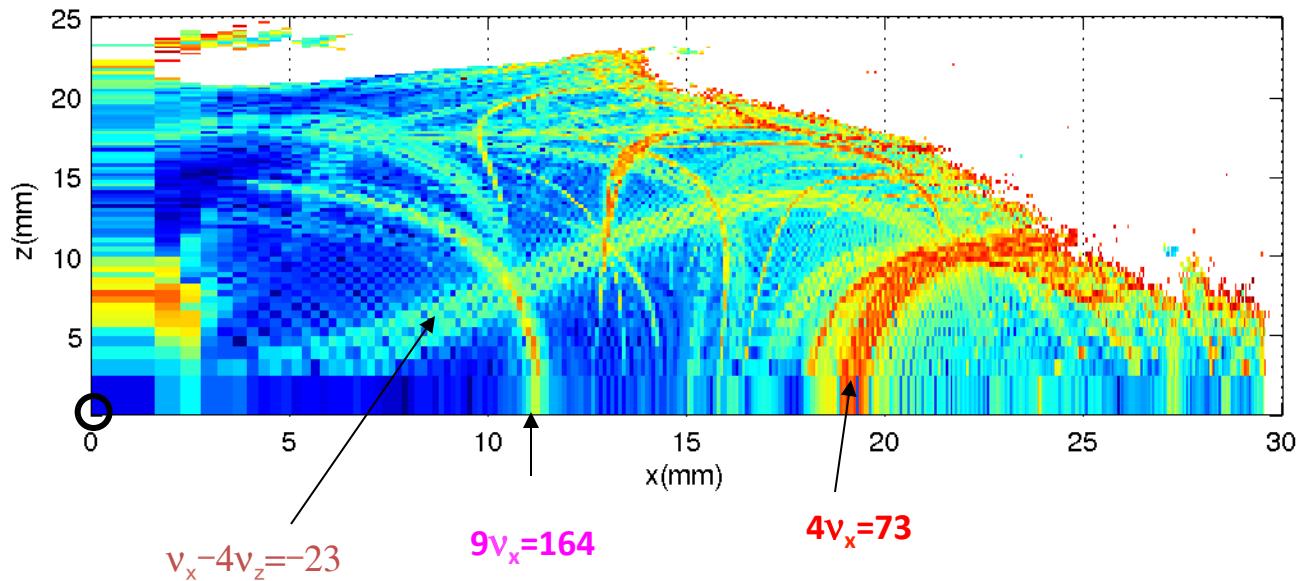
WP sitting on
Resonance node

$$v_x + 6v_z = 80$$

$$5v_x = 91$$

$$v_x - 4v_z = -23$$

$$2v_x + 2v_z = 57$$

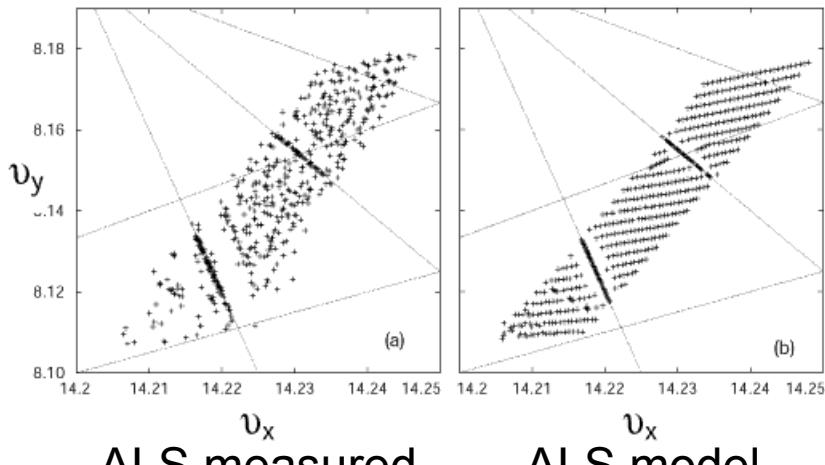


Frequency Map Analysis: ALS and BESSY-II

P. Kuske (BESSY-II)

**ALS linear lattice corrected to
0.5% rms β -beating**

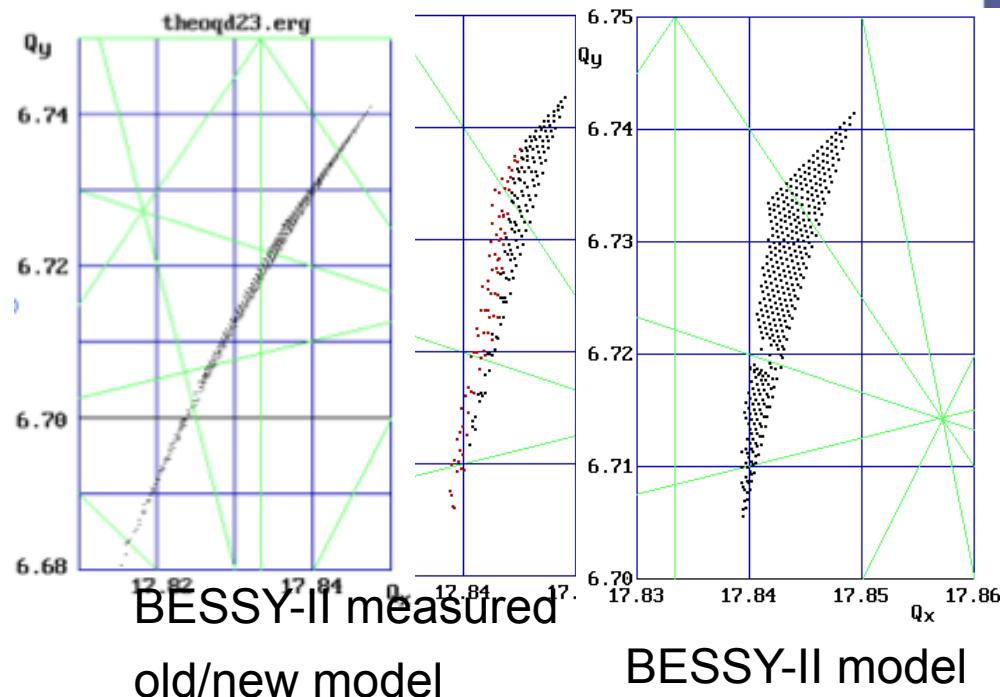
FM computed including residual
 β -beating and coupling errors



D. Robin et al PRL 85, 3 (2000)

A very accurate description of machine model is mandatory

**BESSY-II with harmonic sextupole
magnets, chromaticity, coupling**



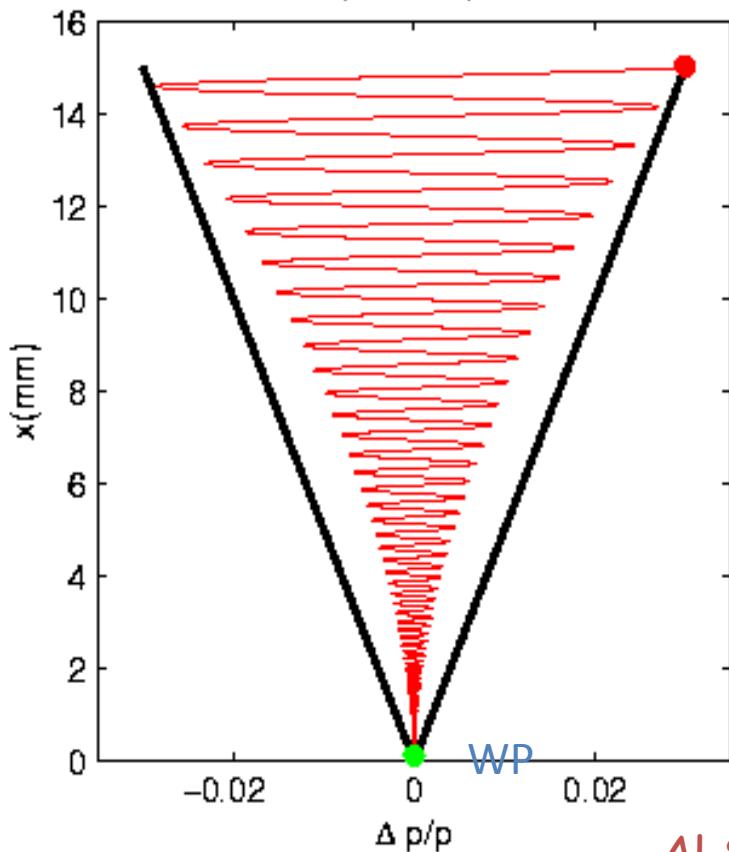
- Fringe fields: dipole, quadrupole (and sextupole) magnets
- Systematic octupole components in quadrupole magnets
- Decapoles, skew decapoles and octupoles in sextupole magnets

Particle behavior after Touschek scattering

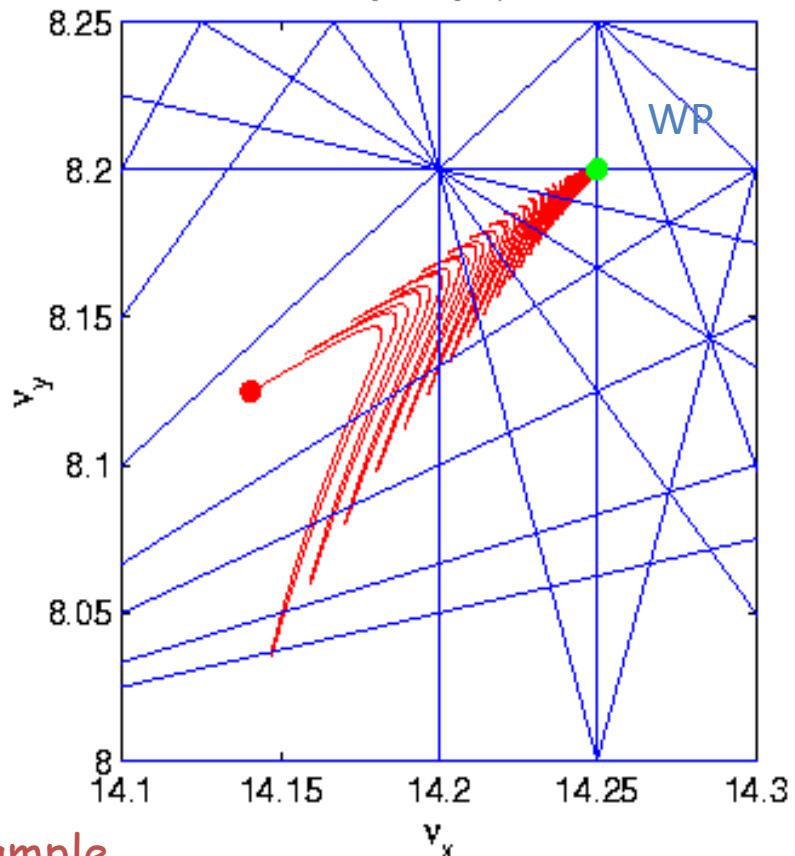
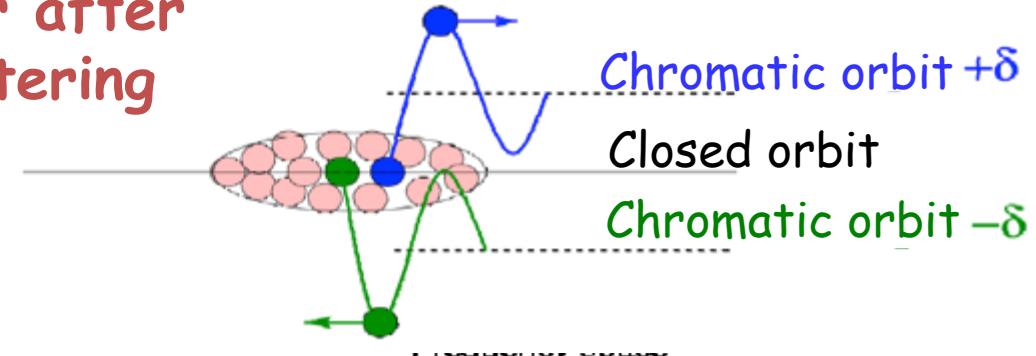
$$x = \sqrt{A_x \beta_{x_1}} + \eta_0 \delta$$

$$A_x = \gamma_{x_0} (\eta_0 \delta)^2 + 2\alpha_{x0} (\eta_0 \delta) (\dot{\eta}_0 \delta) + \beta_{x0} (\dot{\eta}_0 \delta)^2$$

Amplitude space

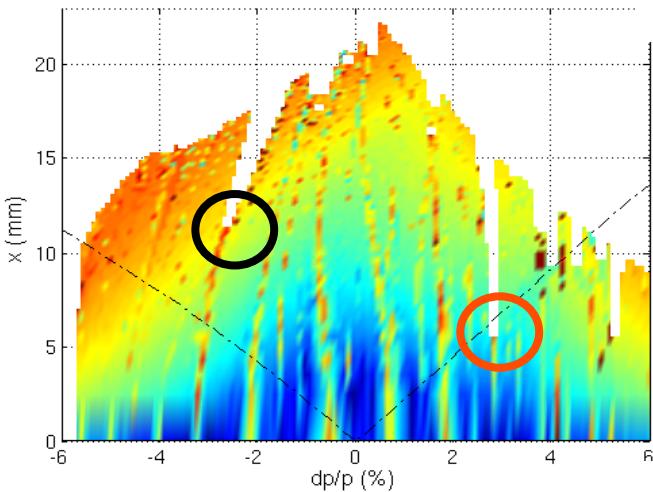


ALS Example

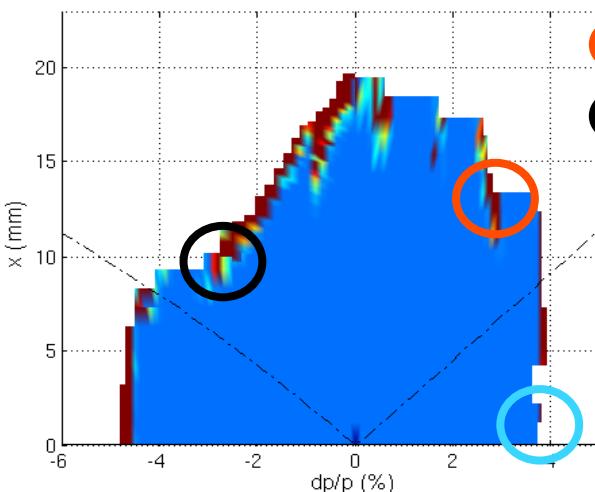
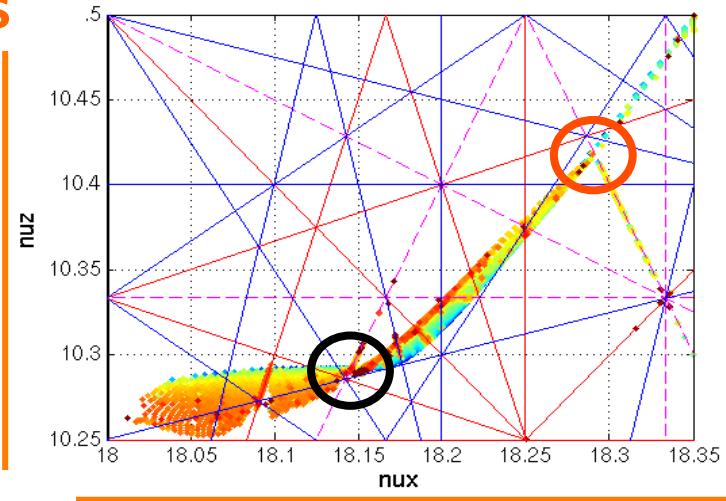


Large off-momentum DA to keep the Touschek scattered particule in its way back to WP

The SOLEIL energy acceptance of the bare machine is large : +/- 4%.
Agreement of a few percents (a factor 2 common 10 year ago)
 Complete optimized linear and non-linear model



Simulations

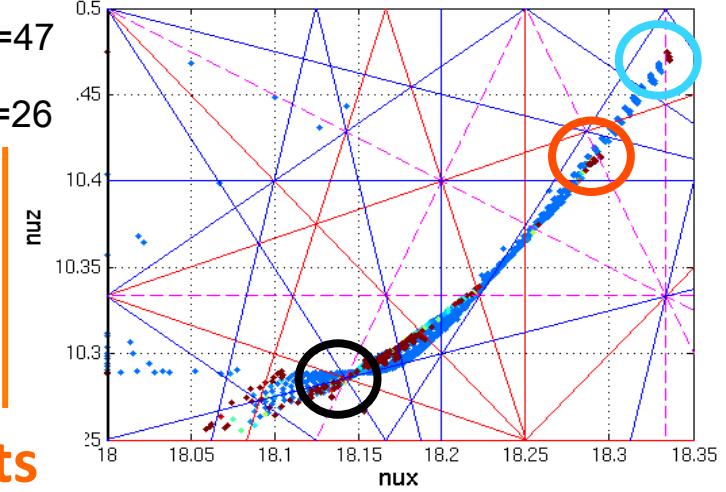


○ $dp/p = + 4 \% : 3v_x=55.$

○ $dp/p = + 3 \% : 2v_x+v_z=47$

○ $dp/p = - 3 \% : 2v_x-v_z=26$

Measurements



TOUSCHEK LIFETIME LOCAL MOMENTUM ACCEPTANCE

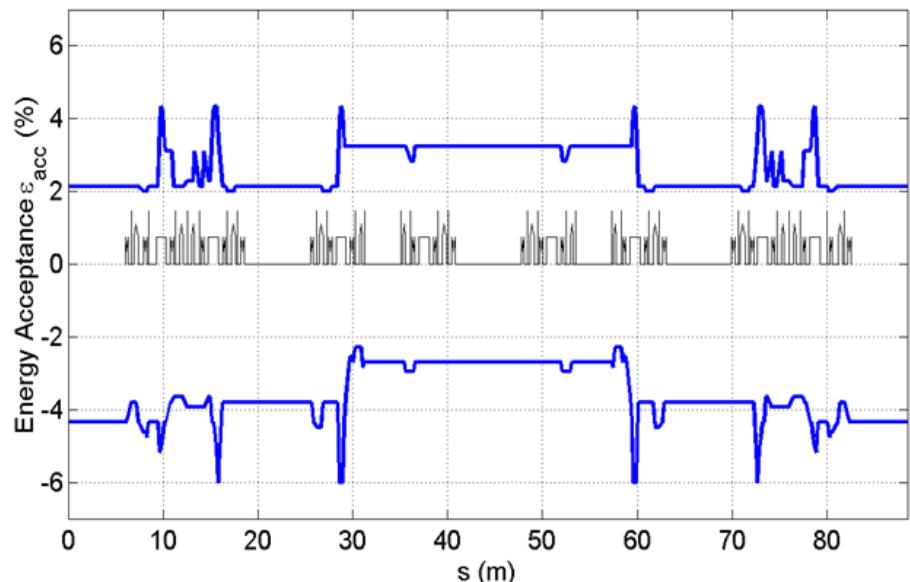
$$\frac{1}{\tau_{T_{1/2}}} = \left(\frac{r_e^2 c N}{8\pi \gamma^3 \sigma_l} \right) \cdot \frac{1}{L} \int_0^L \sigma_x(s) \sigma_z(s) \sigma'_x(s) \varepsilon_{acc}^2(s) ds$$

SOLEIL

Local $\varepsilon_{acc}(s)$ is determined by:

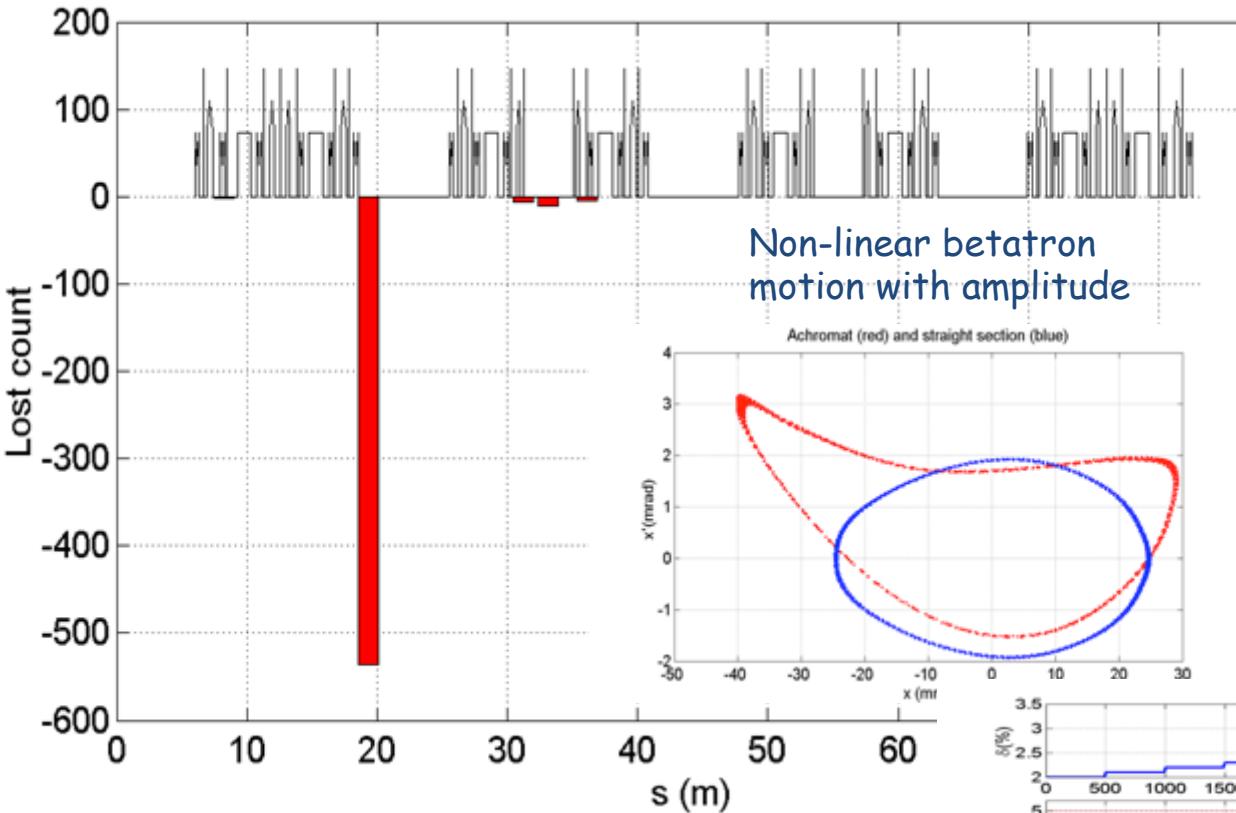
- RF bucket momentum height (longitudinal energy acceptance)
- aperture of the vacuum chamber, or **by dynamic aperture (transverse)**
(if the induced amplitude after a Touschek scattering exceeds one of these two transverse limits).

Working point 18.30 & 10.27
TRACY: 6D tracking w/ 1% coupling

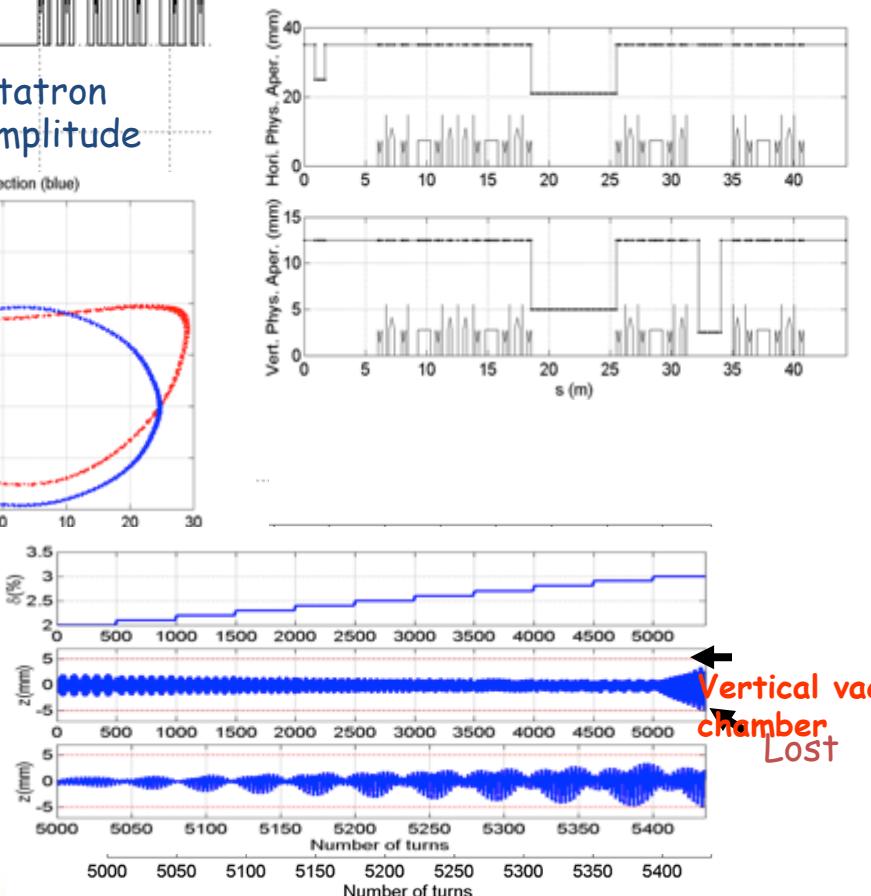


Beam loss location understanding to reducing activation

SOLEIL: Working point 18.30 & 10.27 ($\Delta v=0.03$)



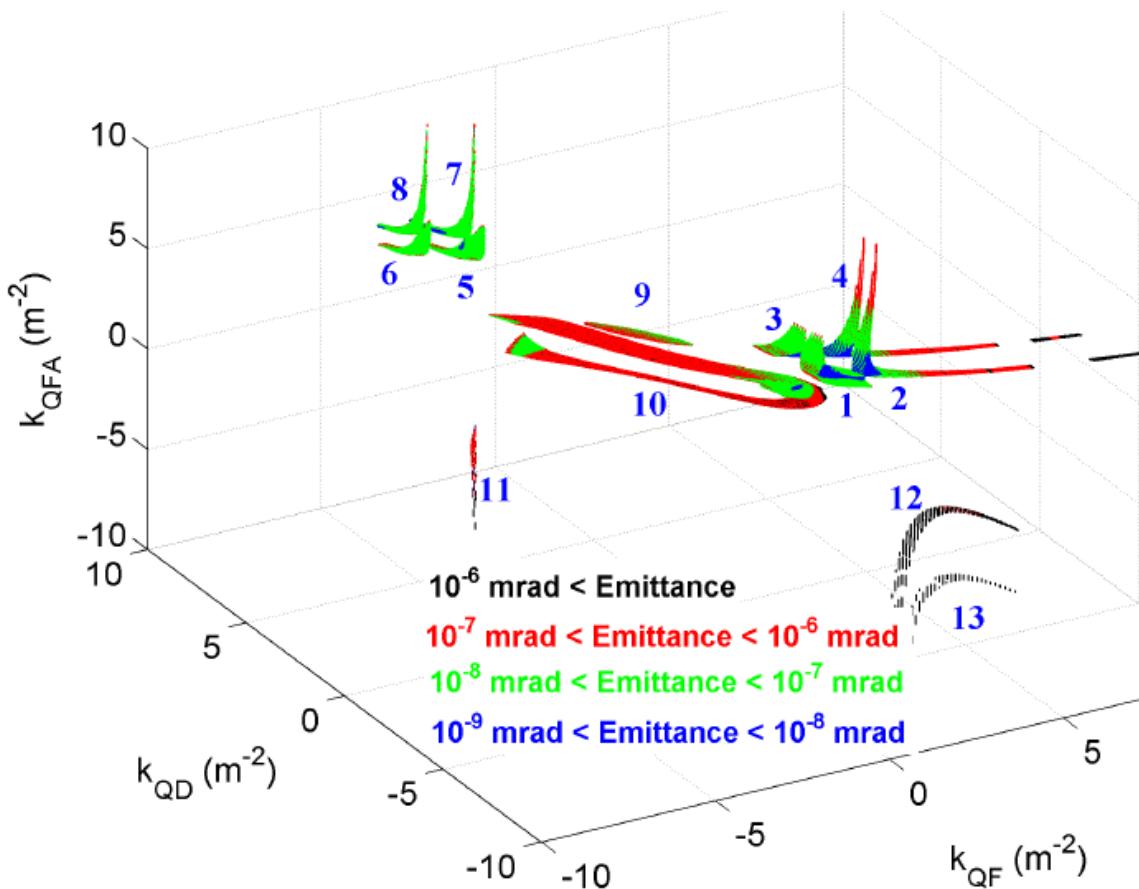
Particles lost in the vertical plane



Lattice optimization

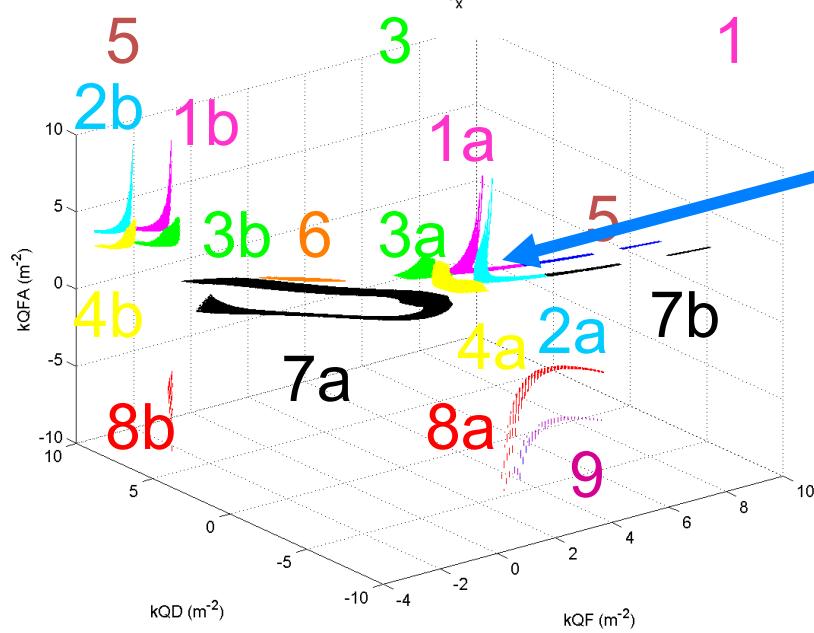
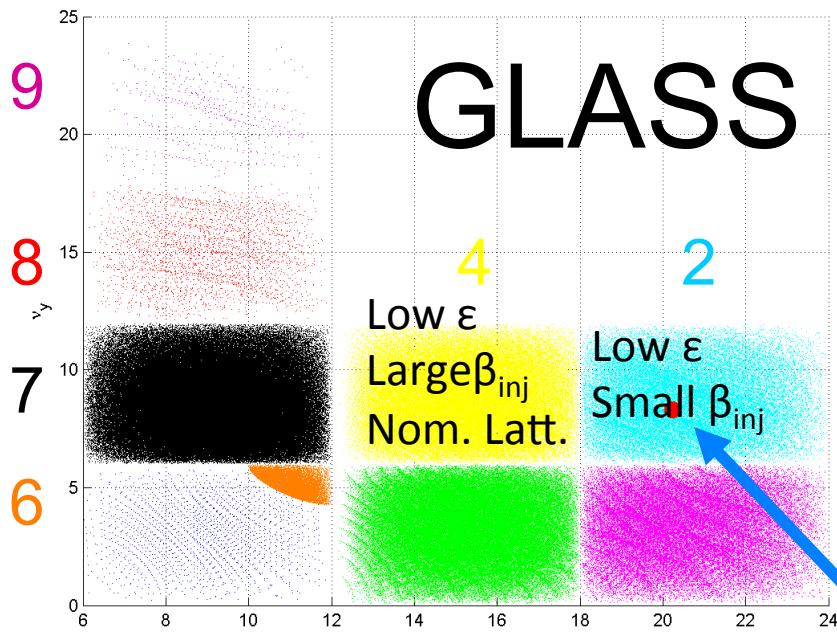
- **GLASS – Global Analysis of All Stable Solutions**
 - Scan for optimum lattice solution for highly periodic lattices (few parameters)
- D. Robin, et al., Physical Review Special Topics 024002 (2008)
- A billion of lattices scanned with 3 quadrupole and 2 sextupole families
 - After 1 day of computation, 1 million of stable solutions
 - Then compute main properties of these solutions to build up a large **exhaustive database**
 - Solutions sorted by emittance values, tunes, DA sizes, momentum apertures but also brilliance
 - **Give a global view of the lattice, very practical**

13 areas of solutions identified counter intuitive and/or promising

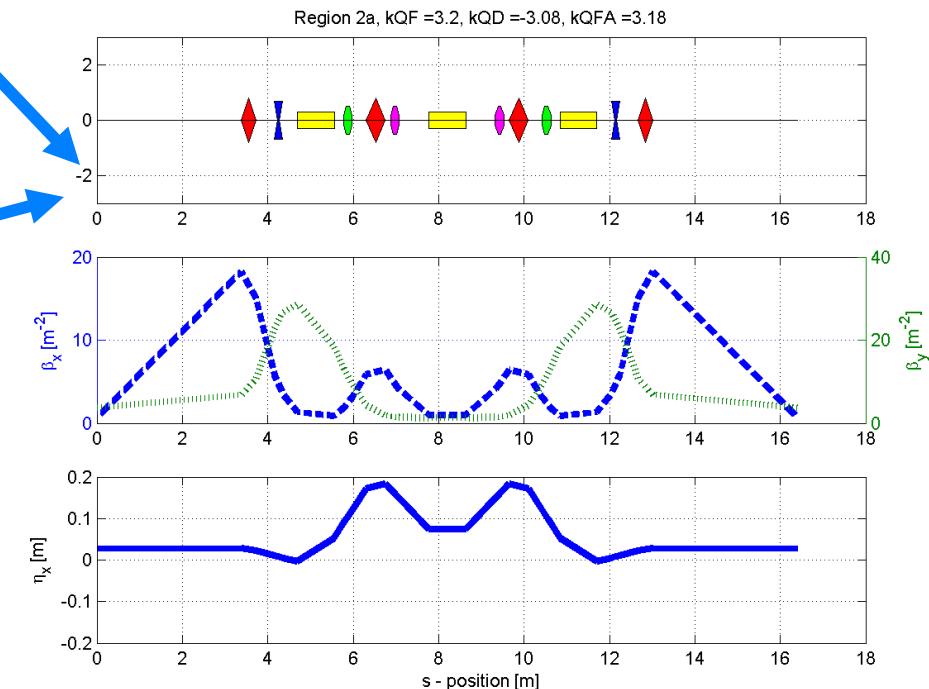


7 Regions with low (<10 nmrad) emittances

Courtesy of D. Robin,
NLBD workshop, Diamond 09



- Low emittance
- Low momentum compaction
- Small beta functions in center bend
- Small horizontal beta in straights



Courtesy of C. Steier

Multi-Objective Genetics Algorithms MOGA

- Exhaustive and global scanning is not possible in a finite time for most 3GLSs
- Indeed large number of parameters (~ 10 families of quadrupoles, of sextupoles)
- **Genetic algorithms** are a very promising solution
 1. Based on direct tracking
 2. Open new optimization windows
 3. Give solutions never thought about
 4. Beam-based checked
- Work started since almost 10 years at APS and followed by other labs and starts to give nice and practical results
- APS (M. Borland et al), ALS, BNL, LSLN, TPS, ...

Merit functions

- NSLS II: Size of on ($x > 15$ mm) / off-momentum dynamics aperture (2-3 %)

$$f_1 = S(\delta = 0) \quad f_2 = S(\delta = -2.5\%) + S(\delta = 2.5\%),$$

- Adding tune shift with amplitude: faster convergence

$$f_1 = \sum_{\delta} S(\delta, y = 1\mu\text{m})$$

$$f_2 = \left(\frac{\partial \nu_x}{\partial J_x} \right)^2 + \left(\frac{\partial \nu_x}{\partial J_y} \right)^2 + \left(\frac{\partial \nu_y}{\partial J_y} \right)^2,$$

- NSLSII: 3 harmonics + 6 chromatic sextupoles
- Tracking codes: Pelegant (APS), Tesla (BNL), ...
- Small number of turns for DA tracking (64-128)
- Then results post filtered using FMA to select best and robust solutions

Maximal survival area in δ -x plane

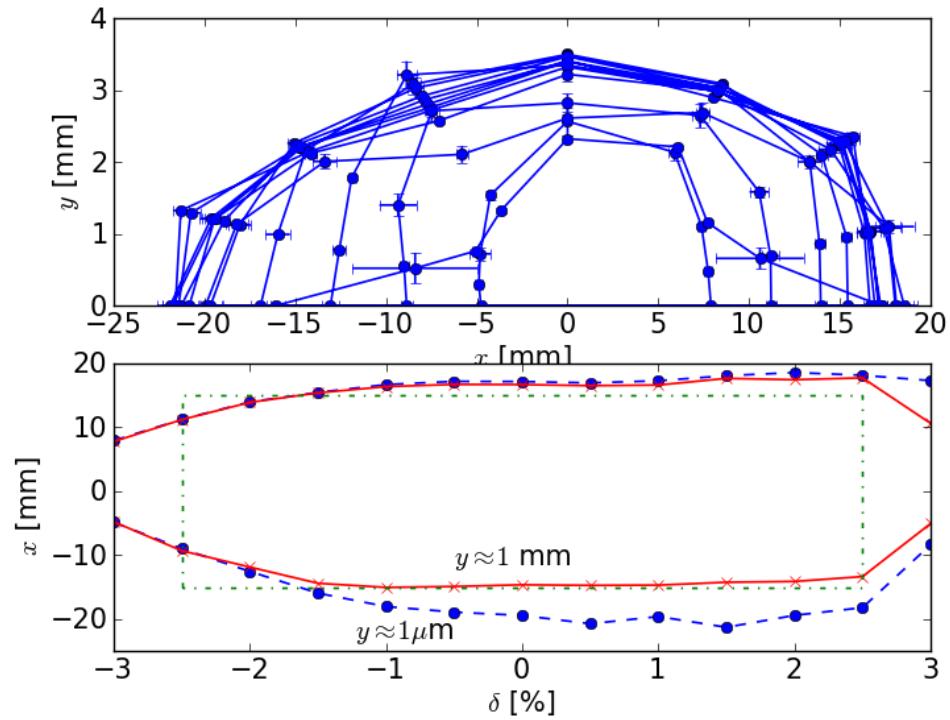
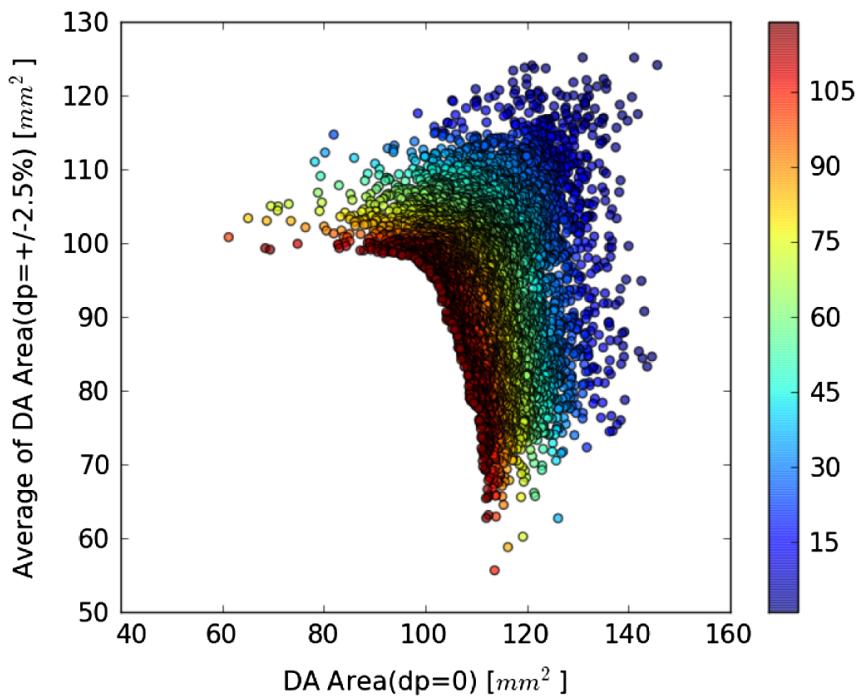


FIG. 1. The last generation of objective functions: DA of on-momentum (horizontal axis) and off-momentum (vertical axis) particles. Points are colored according to their rank.

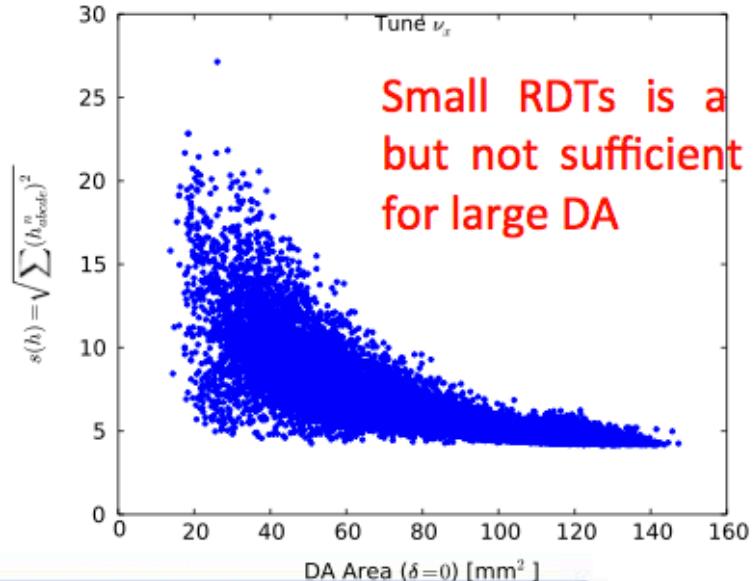
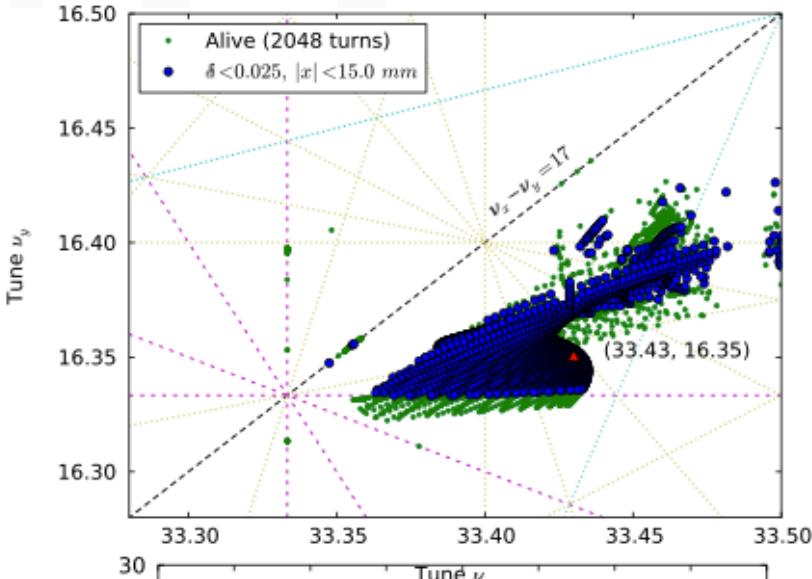
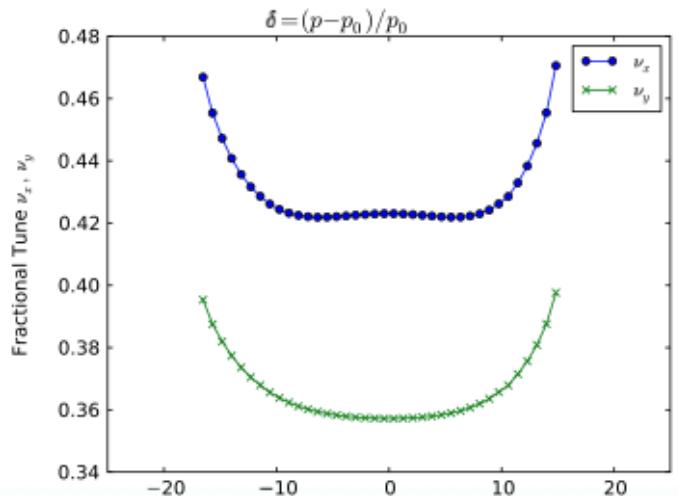
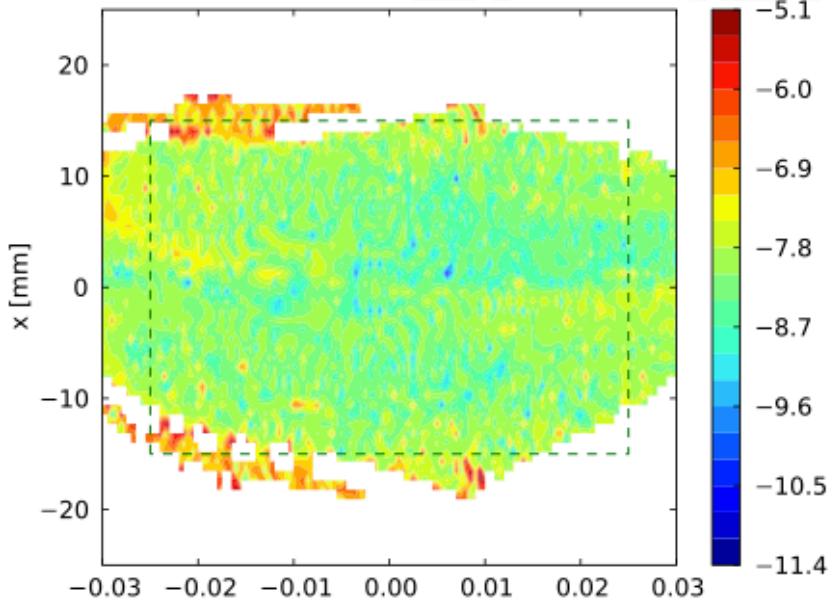
Population of 6000, computation of 300 generation take less than a week on 96 Xeon 2.33 GHz CPU in a Sun Grid Engine cluster. **L Yang, PRSPAB 14 054001 (2011)**



Multiobjective optimization of dynamic aperture

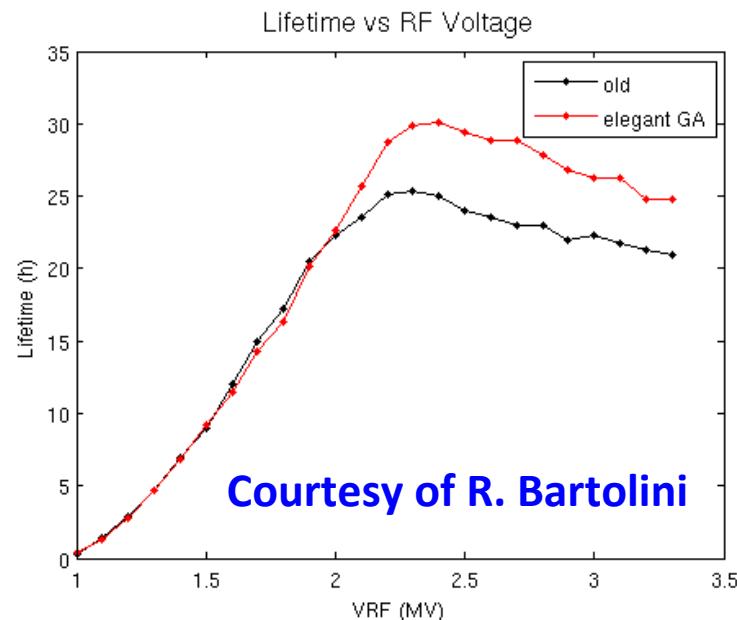
Lingyun Yang, Yongjun Li, Weiming Guo, and Samuel Krinsky

Photon Sciences Directorate, Brookhaven National Laboratory, Upton, New York 11973, USA



Amazing results verified by experiments

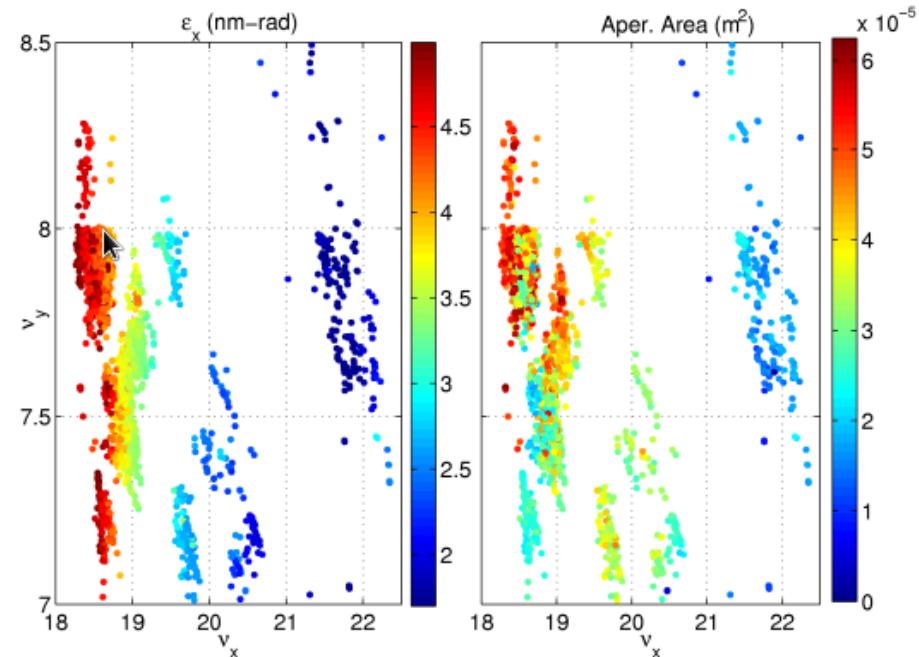
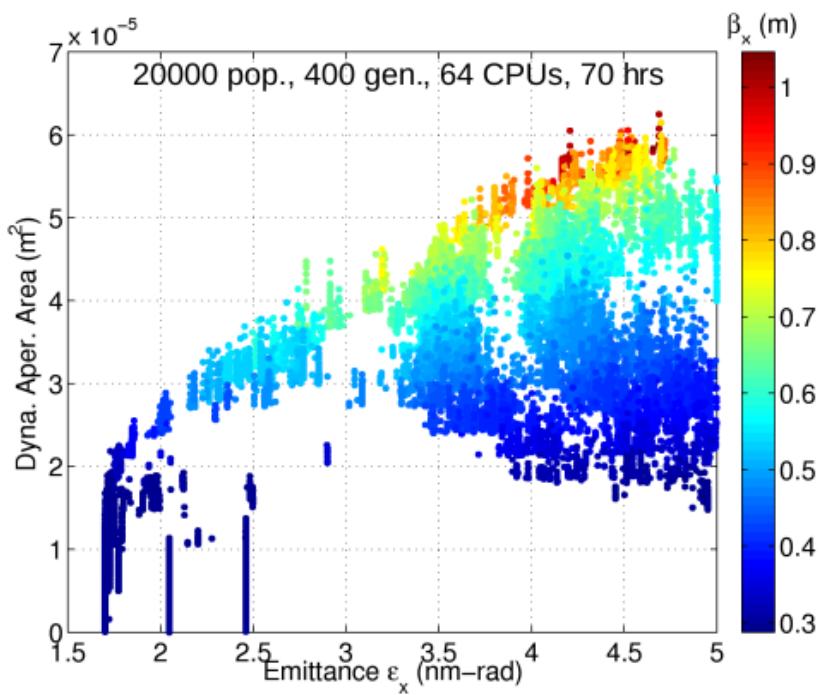
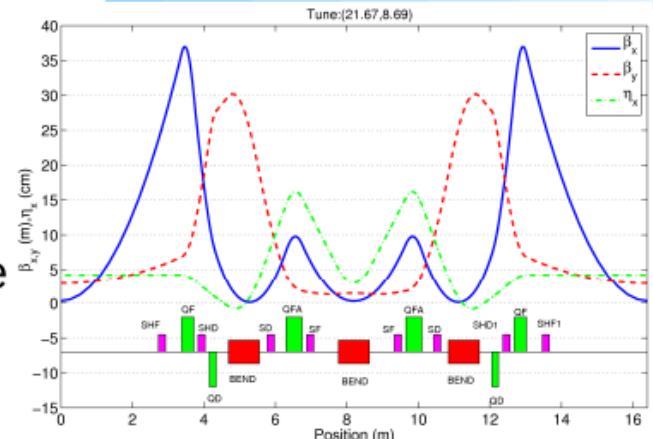
- Multi-objective optimization is effective, pretty robust to errors and very promising
- Experiment results at APS ([M. Borland et al., PAC09](#))
 - Optimization of operation model
 - 24-bunch filling pattern (chro. 6): lifetime improved by 25%
 - Hybrid mode (chro. 11): lifetime improved by 10%
 - Breaking the symmetry of the sextupole: Lifetime 25% better APS slightly than operational lattice as predicted with the same injection efficiency (90-100%)
 - Application to Diamond Light Source: 25 % lifetime improvement.



Simultaneous Linear and Nonlinear Optimization

ALS ultimate lattice is used as an example:

- **7 parameters:** 3 Quads + 6 Sextupoles
- **constraints:** stability, positive partition number, reasonable optics functions
- **3 objectives:** emittance, betax and dynamics aperture



The “standard method” for sextupole optimization

J. Bengtsson, *The sextupole scheme for the SLS: an analytic approach*,
 Internal report SLS-TME-TA-1997-12

a) get the sextupole [+quadrupole] Hamiltonian:

$$\implies \int_{\text{cell}} [H_2(s) + H_3(s)] ds = \sum h_{jklmp} \text{ with}$$

$$h_{jklmp} \propto \sum_n^{N_{\text{sext}}} (b_3 L)_n \beta_{xn}^{\frac{j+k}{2}} \beta_{yn}^{\frac{l+m}{2}} D_n^p e^{i\{(j-k)\phi_{xn} + (l-m)\phi_{yn}\}} - \left[\sum_n^{N_{\text{quad}}} (b_2 L)_n \beta_{xn}^{\frac{j+k}{2}} \beta_{yn}^{\frac{l+m}{2}} e^{i\{(j-k)\phi_{xn} + (l-m)\phi_{yn}\}} \right]_{p \neq 0}$$

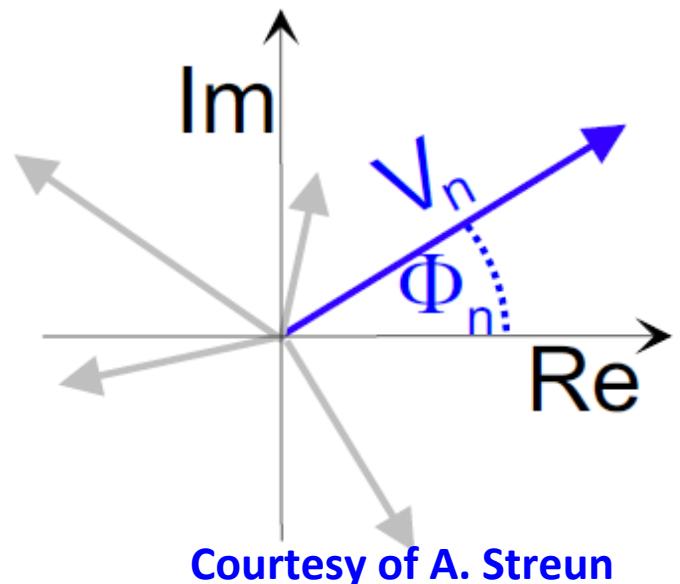
$$h = \sum_n^{N_{\text{sext}}} V_n e^{i\Phi_n} [+ \dots \text{quads for } p \neq 0 \dots]$$

Sextupole_n \leftrightarrow complex vector:

Length $V_n = V_n (b_3, L, \beta_x, \beta_y, D)$

Angle $\Phi_n = \Phi_n (\phi_x + \phi_y)$

- $\Phi_n = 0 \forall n \rightarrow$ tune shifts
- $\Phi_n \neq 0 \rightarrow$ resonances



First order sextupole [+quadrupole] Hamiltonian

- 2 phase independant terms → chromaticities:

$$h_{11001} = +J_x \delta \left[\sum_n^{N_{sext}} (2b_3 L)_n \beta_{xn} D_n - \sum_n^{N_{quad}} (b_2 L)_n \beta_{xn} \right] \rightarrow \xi_x$$

$$h_{00111} = -J_y \delta \left[\sum_n^{N_{sext}} (2b_3 L)_n \beta_{yn} D_n - \sum_n^{N_{quad}} (b_2 L)_n \beta_{yn} \right] \rightarrow \xi_y$$

- 7 phase dependant terms → resonances: $h^N := h$ for N cells, $N \rightarrow \infty \implies$

$$|h_{jklmp}^\infty| = \frac{|h_{jklmp}|}{2 \sin \pi [a_x Q_x^{\text{cell}} + a_y Q_y^{\text{cell}}]}$$

$$a_x = (j - k) \quad a_y = (l - m)$$

$$h_{21000} = h_{12000}^* \longrightarrow Q_x$$

$$h_{30000} = h_{03000}^* \longrightarrow 3Q_x$$

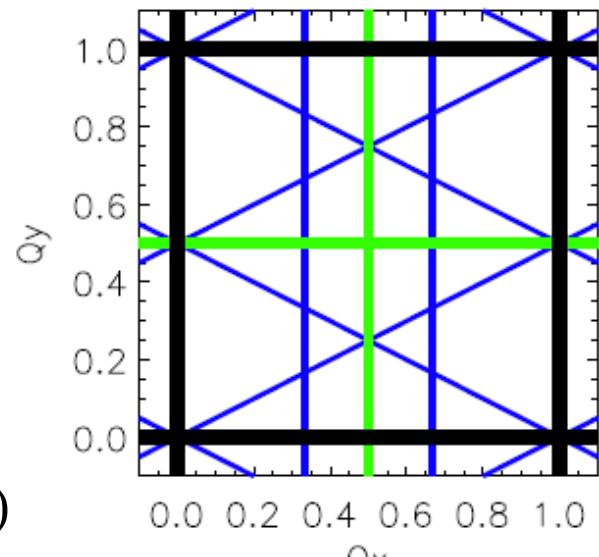
$$h_{10110} = h_{01110}^* \longrightarrow Q_x$$

$$h_{10200} = h_{01020}^* \longrightarrow Q_x + 2Q_y$$

$$h_{10020} = h_{01200}^* \longrightarrow Q_x - 2Q_y$$

$$h_{20001} = h_{02001}^* \longrightarrow 2Q_x \}$$

$$h_{00201} = h_{00021}^* \longrightarrow 2Q_y \} \rightarrow d\beta/d\delta \ (\delta = \Delta E/E)$$



Courtesy of A. Streun

e) ... still not the end: 13 more terms in 2nd order: 5 real, 8 complex
 (Pandora's box has a false bottom!)

Second order sextupole [+first order octupole] Hamiltonian

$$\sum_n \sum_m (b_3 L)_n (b_3 L)_m \times (\beta_n, \phi_n \beta_m, \phi_m \dots) + \left[\sum_q (b_4 L)_q \times (\beta_q, \phi_q \dots) \right]$$

- 3 phase independant terms → amplitude dependant tune shifts:

$$\frac{\partial Q_x}{\partial J_x} \quad \frac{\partial Q_x}{\partial J_y} = \frac{\partial Q_y}{\partial J_x} \quad \frac{\partial Q_y}{\partial J_y}$$

- 2 phase independant off-momentum terms → second order chromaticities:

$$\xi_{x/y}^{(2)} = \frac{\partial^2 Q_{x/y}}{\partial \delta^2}$$

- 8 phase dependant terms

→ octupolar resonances:

$$h_{40000} \rightarrow 4Q_x$$

$$h_{31000} \rightarrow 2Q_x$$

$$h_{00400} \rightarrow 4Q_y$$

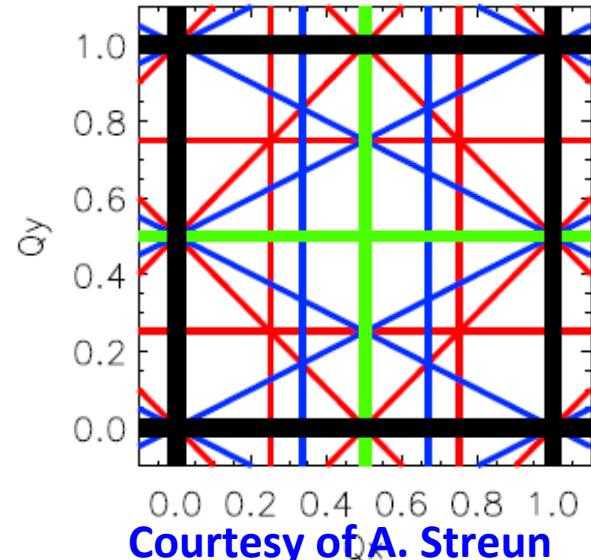
$$h_{20110} \rightarrow 2Q_x$$

$$h_{20200} \rightarrow 2Q_x + 2Q_y$$

$$h_{00310} \rightarrow 2Q_y$$

$$h_{20020} \rightarrow 2Q_x - 2Q_y$$

$$h_{01110} \rightarrow 2Q_y$$



Versatile Sextupoles

all 120 sextupoles were delivered with H&V corrector coils
⇒ make skew quadrupoles and auxiliary sextupoles

120 sextupoles in 9 families:

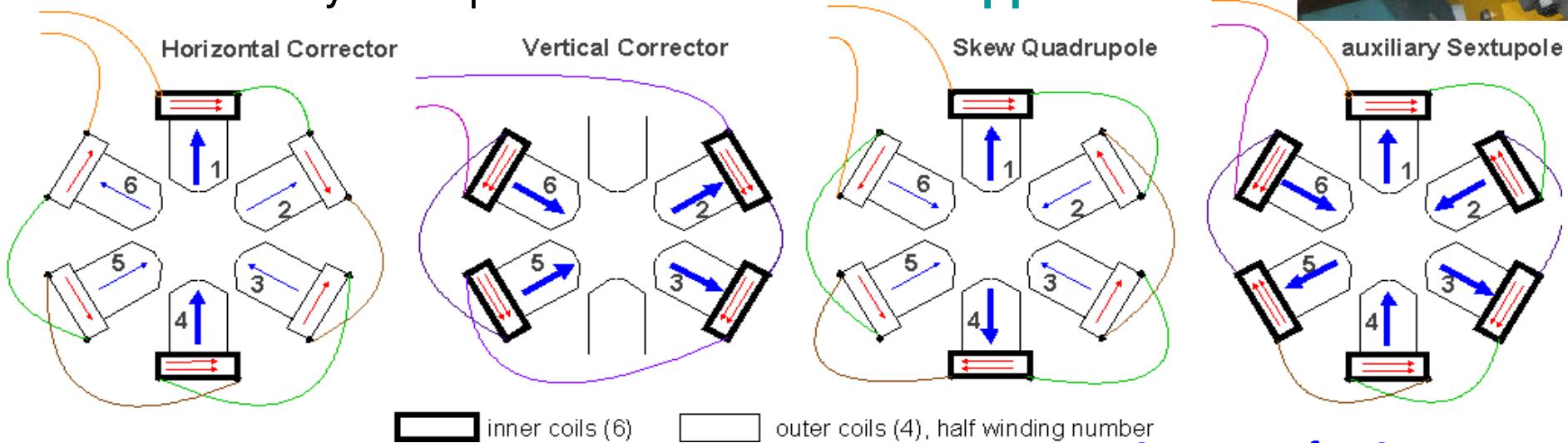
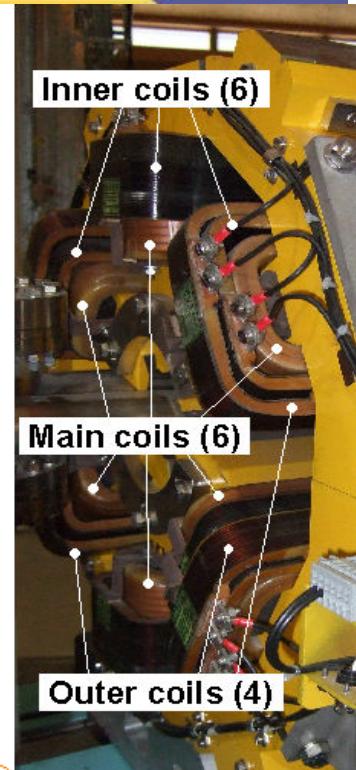
SF(24), SD(24), SE(24) → **chromaticities**

SSA(12), SSB(12), SMA(6), SMB(6), SLA(6), SLB(6) → **D.A.**

SD, SE, S*B: **72** H&V correctors → **orbit correction**

S*A: **24** skew quads ($\eta=0$) → **betatron coupling**

SF: **12** skew quads ($\eta>0$) → **vertical dispersion**
12 auxiliary sextupoles → **resonance suppression**



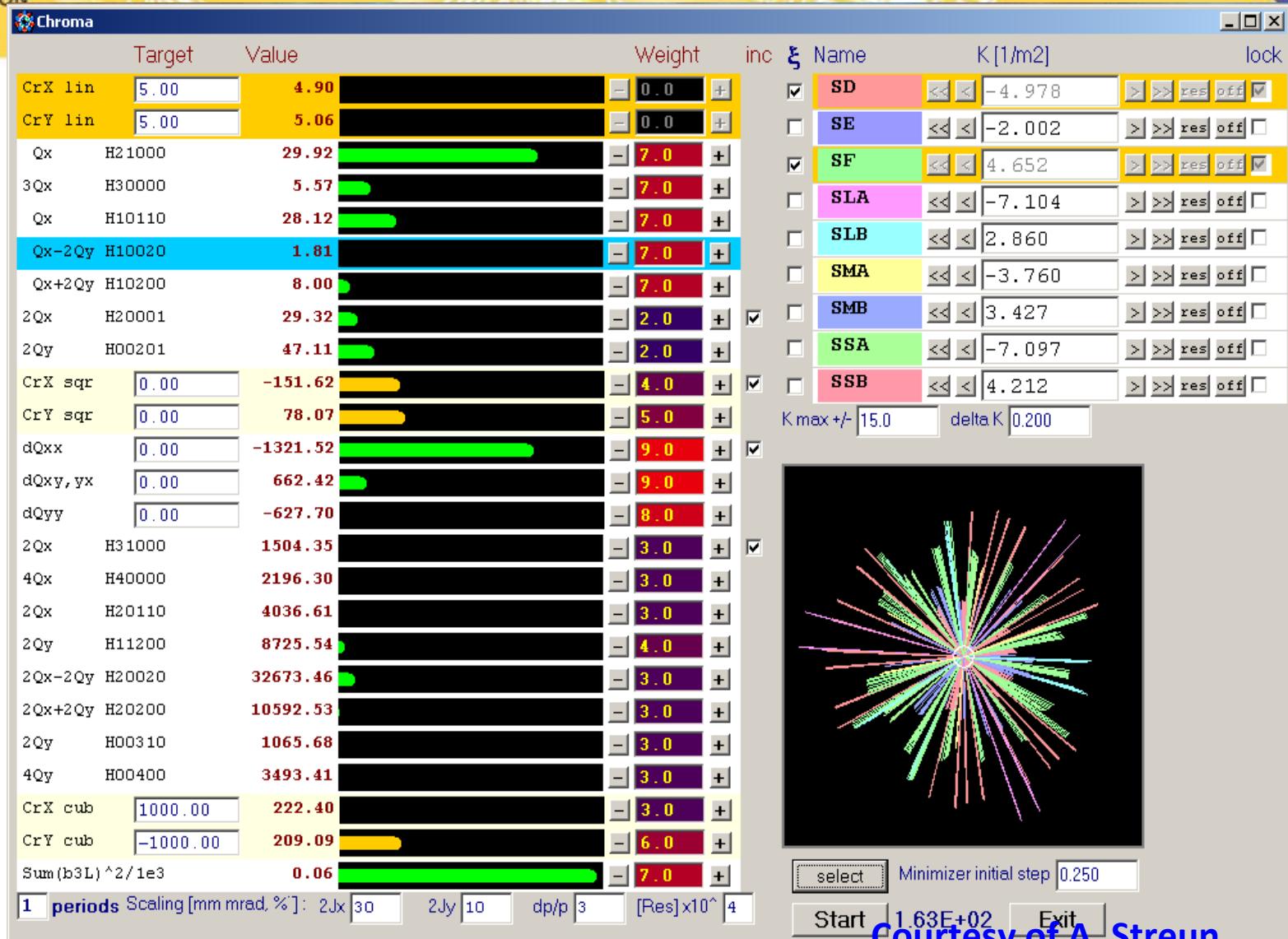
Courtesy of A. Streun

f) Tool for sextupole optimization (OPA)

Analytical
expressions
for 1st and 2nd
Hamiltonian
modes.
(J.Bengtsson)

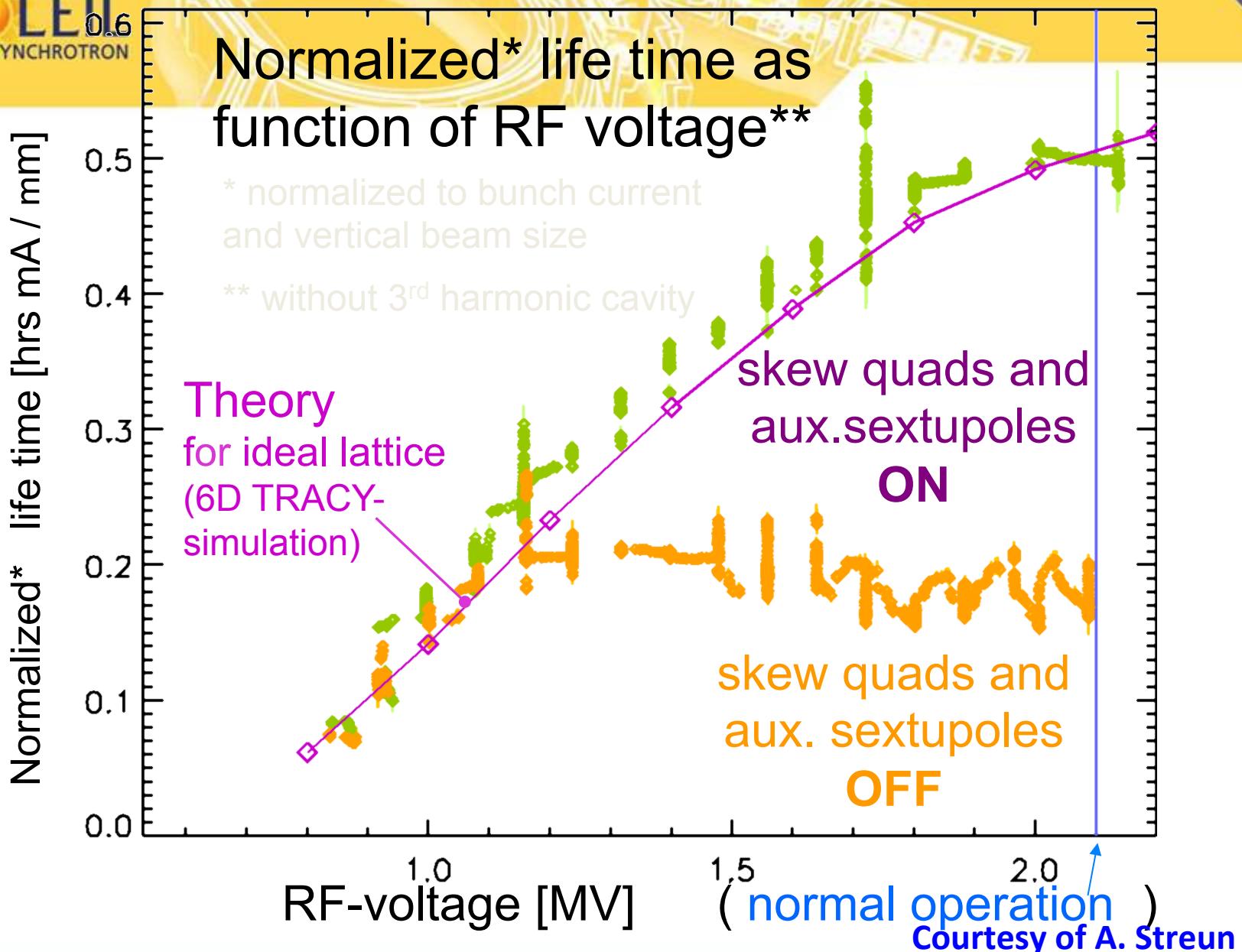
Numeric
differentiation
for 1st, 2nd, 3rd
chromaticity

$S(b_3 l)^2$
included in
minimization



Courtesy of A. Streun

Lifetime in agreement with design

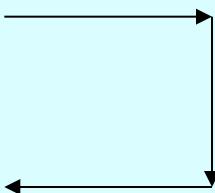


Comparison real lattice to model linear and nonlinear optics

Frequency Maps and amplitudes and phases of the spectral line of the betatron motion can be used to compare and correct the real accelerator with the model

Closed Orbit Response Matrix
from model

Closed Orbit Response Matrix
measured

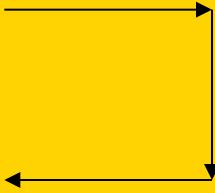


LOCO

Closed orbit based

Spectral lines + FMA
from model

Spectral Lines + FMA
measured



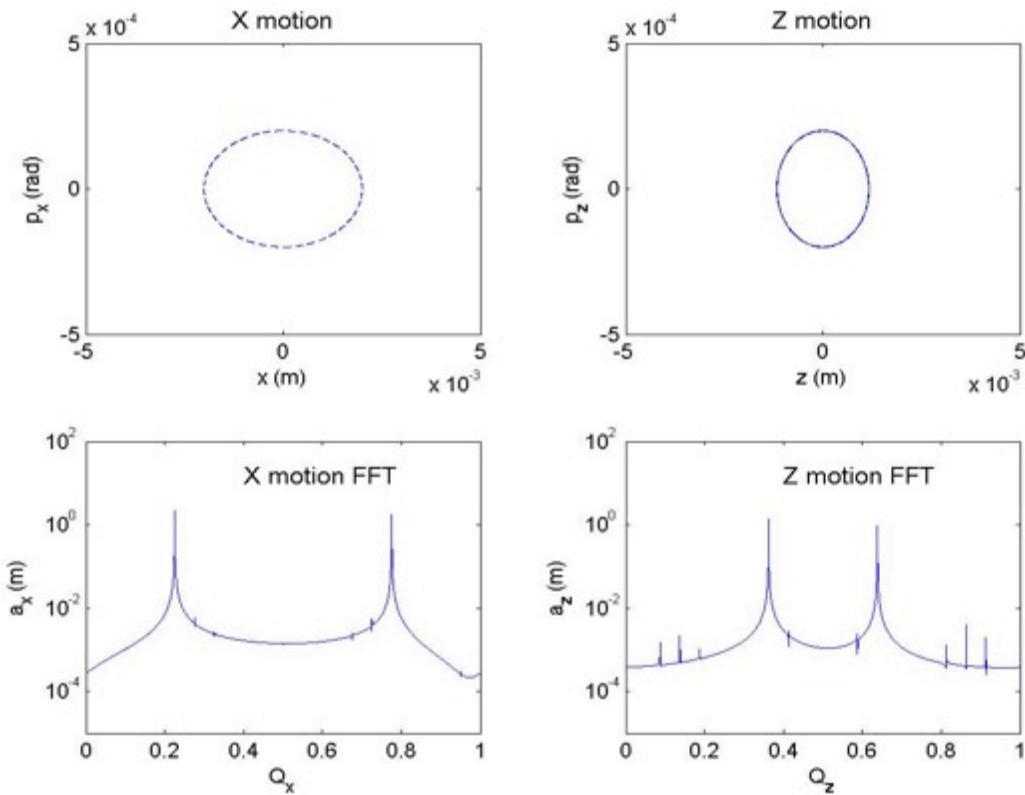
R. Bartolini and F. Schmidt in PAC05

Turn by turn based

Combining the complementary information from FM and spectral line should allow the calibration of the nonlinear model and a full control of the nonlinear resonances

Frequency Analysis of betatron motion

Example: Spectral Lines for tracking data for the Diamond lattice



Spectral Lines detected with SUSSIX (NAFF algorithm)

e.g. in the horizontal plane:

- (1, 0) $1.10 \cdot 10^{-3}$ horizontal tune
- (0, 2) $1.04 \cdot 10^{-6}$ $Q_x + 2 Q_z$
- (-3, 0) $2.21 \cdot 10^{-7}$ $4 Q_x$
- (-1, 2) $1.31 \cdot 10^{-7}$ $2 Q_x + 2 Q_z$
- (-2, 0) $9.90 \cdot 10^{-8}$ $3 Q_x$
- (-1, 4) $2.08 \cdot 10^{-8}$ $2 Q_x + 4 Q_z$

Each spectral line can be associated to a resonance driving term

J. Bengtsson (1988): CERN 88-04, (1988).

R. Bartolini, F. Schmidt (1998), Part. Acc., **59**, 93, (1998).

R. Tomas, PhD Thesis (2003)

Courtesy of R. Bartolini

Frequency Analysis of Betatron Motion and Lattice Model Reconstruction

Using the measured amplitudes and phases of the spectral lines of the betatron motion we can build a fit procedure to calibrate the nonlinear model of the ring

Accelerator Model



- tracking data at all BPMs
- spectral lines from model (NAFF)

Accelerator



- beam data at all BPMs
- spectral lines from BPMs signals (NAFF)

e.g. targeting more than one line

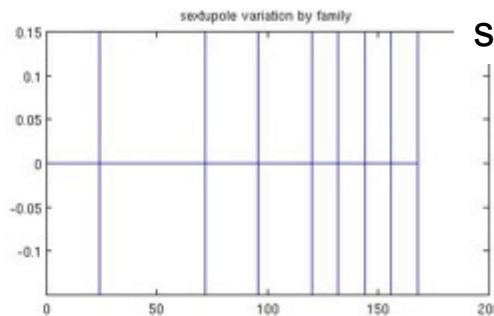
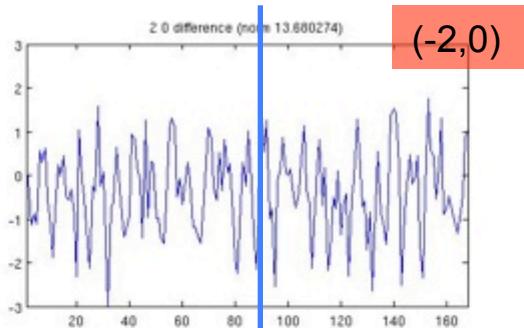
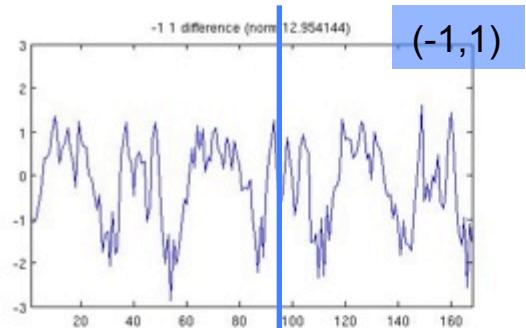
$$\bar{A} = \left(a_1^{(1)} \quad \dots \quad a_{NBPM}^{(1)} \quad \varphi_1^{(1)} \quad \dots \quad \varphi_{NBPM}^{(1)} \quad a_1^{(2)} \quad \dots \quad a_{NBPM}^{(2)} \quad \varphi_1^{(2)} \quad \dots \quad \varphi_{NBPM}^{(2)} \quad \dots \right)$$

Define the distance between the two vector of Fourier coefficients

$$\chi^2 = \sum_j (A_{Model}(j) - A_{Measured}(j))^2$$

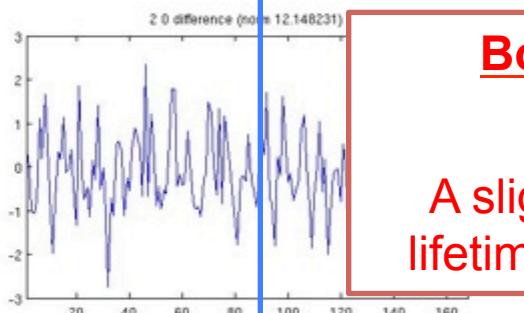
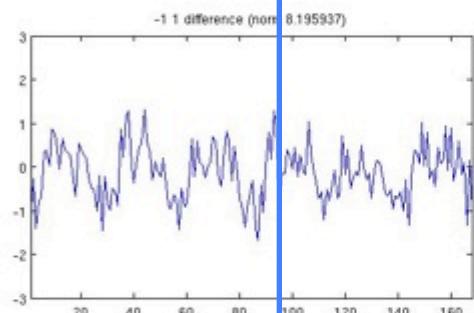
Least Square Fit of the sextupole gradients to minimize the distance χ^2 of the two Fourier coefficients vectors

Simultaneous fit of (-2,0) in H and (1,-1) in V



sextupoles

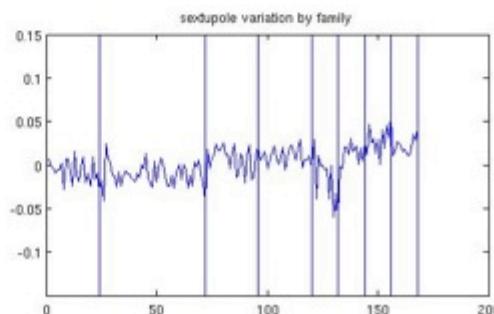
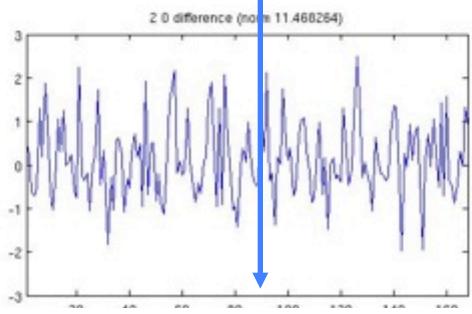
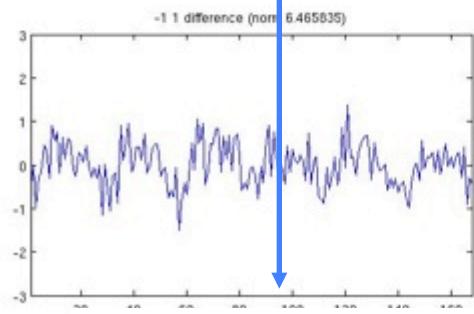
start



Both resonances are controlled

iteration 1

A slight improvement in the lifetime (10%) was measured



iteration 2

Both resonance driving terms are decreasing

Courtesy of R. Bartolini

Conclusion (1)

- Tracking codes trustable, long term tracking, parallelized
- Agreement between model and experience: tune shift with amplitude, DA, non-linear chromaticities are good providing a good linear and non-linear model
- Model efficiently used for predicting performance, impact of IDs using turn by turn data
- FMA like techniques give a global view of dynamics
 - Shift from 1D to 2D view
- Promising use RDTs based on TbT data (coupling correction¹, Non-linear LOCO like): will benefit from higher resolution TbT data
- MOGA enables us to optimize on/off dynamics aperture, use DAs, FMAs, RDTs in merit functions.
 - Shift from sequential to simultaneous

¹ talk by A. Franchi this afternoon

Conclusions (2)

- **3 GLSs run top-up operation mode**
 - MTFB >7 days
 - High injection efficiency (low beam loss rate, green technology, accelerators ...)
- **Preserving high performances**
 - Need to monitor linear and non-linear dynamics (Twiss functions, RDTs, Tunes, Chromaticities, ..)
 - Soon ~100 pm RMS beam stability
 - Strong need for beam-based measurement
 - Beam based alignment, girder based alignment
 - Multi-pole beam based correction + Feedback syst.
 - Beam-based coupling correction + feedback syst.
 - ...

Development of on-line continuous tools to measure beta-beats, chromaticity evolutions, local coupling, performance degradation but without perturbing the user experiments

Perspectives

- **Pushing high performance diagnostics**
 - TbT BPM performance, imaging, emittance measurement, ...
- **Pushing ID optimization**
- **Wanted features**
 - To what extent can beam-based measurement replace magnetic measurements ?
 - RDT measurement A. Franchi, R. Thomas, F. Schmidt, PRSTAB 10, 074001 (2007)
 - SOLEIL: strong limitation for probing off-momentum dynamics ($\delta\text{-}x$ maps). Large RF frequency shifts make aging the cold tuning system of the superconductive cavities
 - A **special cavity for kicking** the beam energy over one turn would be very valuable

Conclusions and Perspectives

Over the years

New Techniques have come out

New tools have emerged

Mass computing has become affordable

New way of optimization

Open new horizon windows

Push limits of Accelerators

Makes us confident for the design and operation of ever brighter light sources, diffracted limited in H&V planes

Acknowledgments

- M. Borland and L. Emery (APS, USA)
- P. Kuske (HZB/Bessy II, Germany)
- X. Huang , Y. Cai, J. Safranek (SLAC, USA)
- C. Steier, D. Robin (ALS, USA)
- L. Yang (BNL, USA)
- K. Sotome (SPring8, Japan)
- S. C. Leemann, MAX-lab, Sweden
- A. Streun (PSI, Switzerland)
- R. Bartolini (Diamond/John Adams Institute, UK)
- A. Franchi (ESRF, France)
- R. Thomas, Y. Papaphillipou (CERN, Switzerland)
- A. Kling (Petra III, Germany)
- M. S. Chiu (TPS, Taiwan)
- SOLEIL Accelerator Physics Group

Mapped IDs à la ESRF

- Nonlinear maps of IDs are generated using the 3D RADIA code.
- To be read in a tracking code.
- BETA-ESRF does not take into account chromatic terms due to the edges of bending magnet.
- BETA-SOLEIL and TRACYII have been modified in order to read the IDs maps.
- Very good agreement between the two codes for ON and OFF momentum dynamic apertures.

- The angular kicks experienced by the particle are derived from the function:

$$\phi(x, z, s) = \left(\int_{-\infty}^s B_x ds' \right)^2 + \left(\int_{-\infty}^s B_z ds' \right)^2$$

- **f** is integrated over 1 period resulting in a potential function **U**:

$$U(x, z) = \int_{\text{1 period}} \phi(x, z, s) ds$$

- The angular kick experienced by a particle over the undulator period is:

$$\Delta x' = -\frac{1}{2(p/e)^2} \frac{\partial U}{\partial x}(x, z)$$

$$\Delta z' = -\frac{1}{2(p/e)^2} \frac{\partial U}{\partial z}(x, z)$$

P. Elleaume, EPAC'92

ID building strategy

❖ **Tolerances**

$$\int B_x ds = \int B_z ds = \pm 20 G.cm \quad \iint B_x dsds' = \iint B_z dsds' = \pm 1 G.m^2$$

❖ **A 3 step-process using ID builder (O. Chubar)**

1. **Assembly: Module sorting according to magnetic measurements**

→ Minimization of first and second integrals

2. **Shimming: using a merit function**

→ Minimization at different gap and phase values with weight factor

i. On axis first & second integral (angle & position) in H & V plane

ii. Skew and normal gradient for new IDs

iii. Phase error < 0.2°

3. **Magic fingers (different gap and phase values with weight factor)**

→ Reduction of high field integral for large transverse amplitudes

❖ **Expected or unexpected effects depending on gap, phase, current values**

❖ Orbit distortion (Feedforward)

❖ Tune, chromaticity, coupling variations

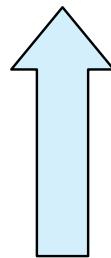
❖ Injection efficiency, lifetime variation (non-linearities, ...)

O.Rudenko and O.Chubar, Proc. of 9th Int. Conf. on PPSN IX, p.362 (2006)

Evaluation:

Ordered
Magnet
Sequence(s)

1 2 3 4 5 6 7 8
5 4 8 1 7 2 6 3



«Decoded»
Undulator
Structure

1 3 5 7
5 4 8 1
7 2 6 3
2 4 6 8

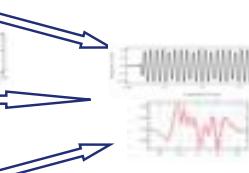
Magnetic Measurements Data
on Individual Magnets (/
Modules) & Partly Assembled
Undulator



Mathematical Model
/ Total Field Calc.
Method



Undulator
Magnetic Field
(/ Field
Integrals)



Characteristics
/ Fitness Terms

Electron Trajectory
Straightness

Radiation Phase
Error

Field Integral
deviation from zero

Integrated Multipoles



Variation Operators for Permutations:

Mutation : - e.g. swap items (magnets) at two randomly chosen positions - [5 4 8 1 7 2 6 3] →

Crossover : - e.g. «order 1» - [1 2 3 4 5 6 7 8] → [??? 4 5 6 7 ?] →
[3 5 6 8 1 2 7 4]

Advantages : object function, arbitrary search space, search from ap population, mutation and

Multi-Objective Genetic Algo.

- 1: Initialize population.
- 2: **repeat**
- 3: crossover: $2 \rightarrow 2$
- 4: mutation: change children.
- 5: calculate obj. func. f_m
- 6: natural selection: “sorting”
- 7: **until** should stop

- ① **Initialize:** the very first generation (random).
- ② **Crossover:** generate children from parents.
- ③ **Mutation:** change the children slightly.
- ④ **Natural selection:** keep population fixed from generation to generation.

Courtesy of L Yang

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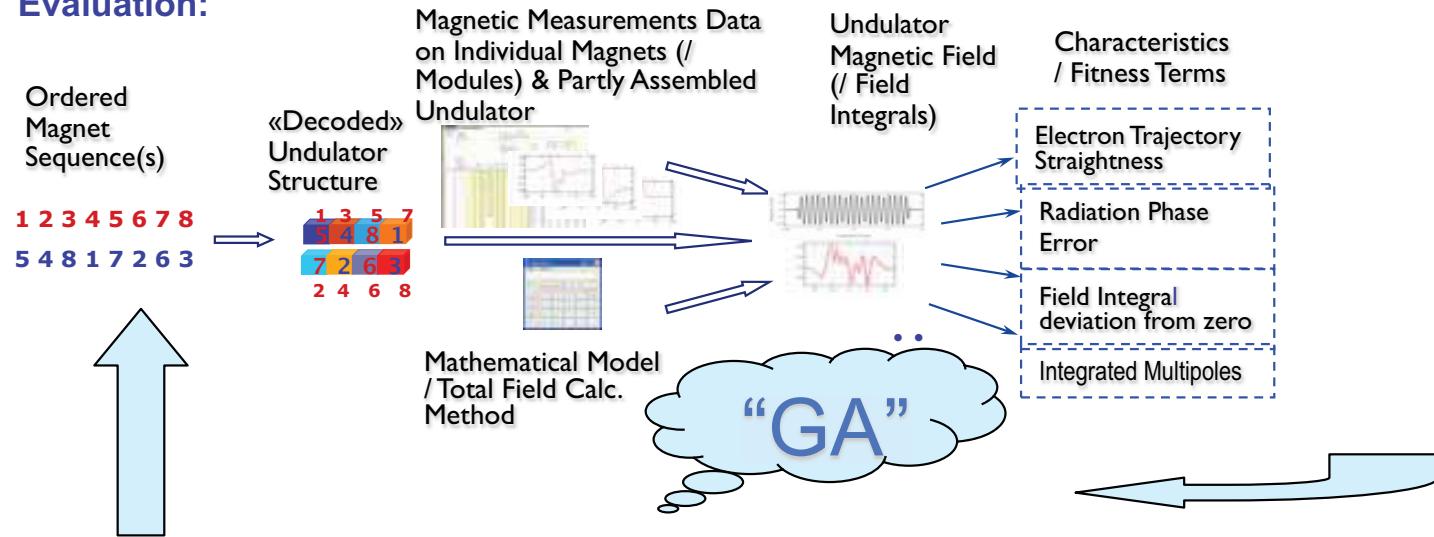
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