DESIGN OF SLOW EXTRACTION SYSTEM AT BOOSTER SYNCHROTRON FOR MUSES

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Abstract

The Booster Synchrotron Ring (BSR) is a part of Multi-USe Experimental Storage rings (MUSES). BSR functions exclusively for the acceleration of ion and electron beams. The maximum accelerating energy is, for example, to be 3 GeV for proton; 1.45 GeV /nucleon for light ions of q/A=1/2; 800 MeV/nucleon for heavy ions of q/A=1/3. Electron beam is accelerated to 2.5 GeV from the injection energy 300 MeV. The accelerated ion and electron beams will be fast extracted and injected into the Double Storage Rings (DSR) by one turn injection. As another operation mode, ion beams will be slowly extracted for the experiments. In this paper, injection and extraction procedures of the BSR, especially slow extraction, are presented.

1 BSR LATTICE DESCRIPTION

As shown in Fig.1, the BSR consists of two arc sections and two long straight sections. Each arc section is mirror symmetrical system, and there are two bending cells. The dispersion in the straight sections outside the arc section is zero. The lattice is specified for eight families of quadrupoles; QF1 and QD1, QF4 and QD4 in the arcs, QF2 and QD2, QF3 and QD3 in the long straight sections.

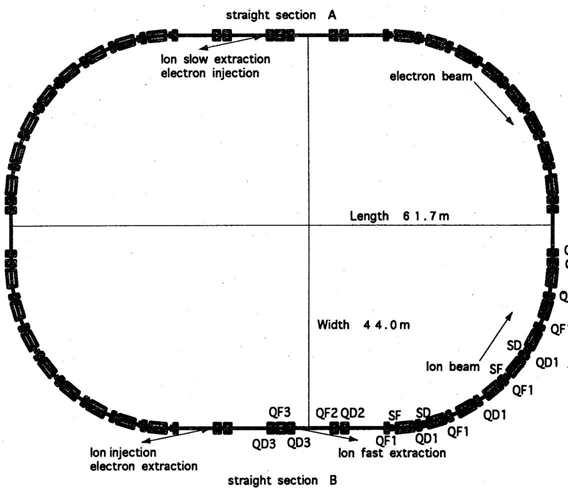

Fig.1 A Layout of BSR

2 INJECTION AND EXTRACTION

An arrangements of the injection and extraction magnets are shown in Figs 2 and 3. The BSR lattice is designed to be able to operate with two different extraction mode for ion beams; fast extraction and slow extraction.

As seen in Fig.2, an electrostatic septum(ES), four septum magnets(SM1, SM2, SM3, SM4) and three bump magnets(BM1) are used for the electron beam multi-turn injection. These magnets except bump magnets, are used

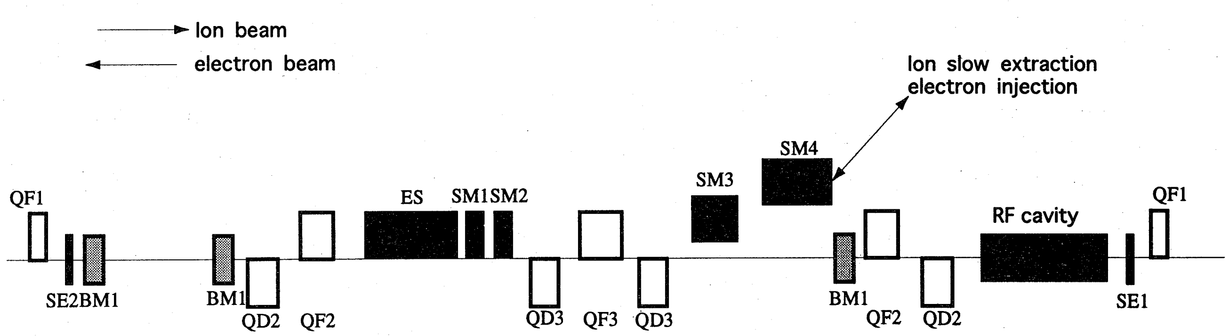

Fig. 2 An arrangement of straight section A

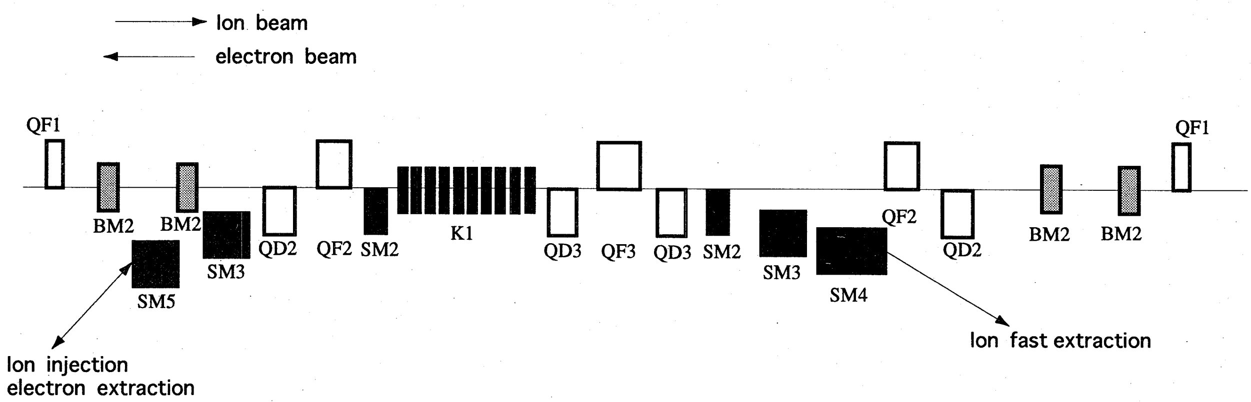

Fig. 3 An arrangement of straight section B

for ion beam slow extraction which is carried out by using third order resonance (ν_{res}=20/3).

As seen in Fig.3, three septum magnets(SM2, SM3, SM5), four bump magnets(BM2) and ten kicker magnets(K1) are used for the ion beam one turn injection. These magnets are also used for electron beam fast extraction. The ion beam from the ACR is injected into the septum magnet SM5 at an angle of with respect to the straight section of the BSR. The ion beam is finally kicked and placed in the reference orbit by the kicker magnets. This kicker magnets are used for ion beam fast extraction. So, kicker magnets must be risen and fallen less than 50 ns.

3 SLOW EXTRACTION

On slow extraction process, at the first, sextupole magnets as a resonance exciter are excited. Secondly horizontal tune is moved from the operating value 6.691 to a value near the third order resonance(20/3) by changing excitation currents of the quadrupole magnets. Beams which have deviated to the distance more than 25 mm inside from the central orbit at the entrance of the ES are deflected inward as large as 3 mrad by the static high voltage of the ES.

The betatron oscillation can be expressed by the following relation

$$\begin{pmatrix} X \\ Y \end{pmatrix} = \begin{pmatrix} \cos\mu-\alpha\sin\mu & \sqrt{1+\alpha^2}\sin\mu \\ -\sqrt{1+\alpha^2}\sin\mu & \cos\mu+\alpha\sin\mu \end{pmatrix} \begin{pmatrix} X_0 \\ Y_0 \end{pmatrix}$$
$$= \begin{pmatrix} 1-\alpha\varepsilon & \sqrt{1+\alpha^2}\,\varepsilon \\ -\sqrt{1+\alpha^2}\,\varepsilon & 1+\alpha\varepsilon \end{pmatrix} \begin{pmatrix} X_0 \\ Y_0 \end{pmatrix}$$

here we put

$$X=\frac{1}{\sqrt{2}}\left[\sqrt{1+\alpha^2}+\alpha\right]^{\frac{1}{2}}(x-y)$$
$$Y=\frac{1}{\sqrt{2}}\left[\sqrt{1+\alpha^2}-\alpha\right]^{\frac{1}{2}}(x+y)$$
$$y=\frac{\beta}{\sqrt{1+\alpha^2}}x' .$$

Let Ψi be the deflection due to the i-th sextupole magnet SEi

$$\Psi_i=g_i x_i^2 ,\ g_i=\frac{B_i''}{2B\rho},$$

where xi is the displacement of the orbit from the equilibrium orbit at SEi and gi is the strength of the Sei. Here we assume that the transfer matrix Γi from electric septum to sextupole magnet is

$$\Gamma_i=\begin{pmatrix} a_i & b_i \\ c_i & d_i \end{pmatrix}.$$

In the third resonance, particles pass each sextupole magnets three times a period. Therefore, taking summations of contributions from each magnet over one period, we have

$$H=-\frac{\varepsilon}{2}\left(X^2+Y^2\right)+\frac{\beta}{3\sqrt{1+\alpha^2}}\Sigma g_i' x_i^2$$
$$=-\frac{\varepsilon}{2}\left(X^2+Y^2\right)+\frac{\beta}{3\sqrt{1+\alpha^2}}\Sigma\left\{\left(X^3-3XY^2\right)g_i'A_i+\left(Y^3-3YX^2\right)g_i'B_i\right\}$$

where

$$g_i'=\frac{\beta_i}{\sqrt{1+\alpha_i^2}}\Big/\frac{\beta}{\sqrt{1+\alpha^2}}g_i ,$$

$$A_i=\frac{3}{8\sqrt{2}}\left(\sqrt{1+\alpha^2}-\alpha\right)^{\frac{3}{2}}\left\{\left(a_i-b_i\right)^3-3\left(\sqrt{1+\alpha^2}+\alpha\right)^2\left(a_i+b_i\right)^2\left(a_i-b_i\right)\right\}$$
$$B_i=\frac{3}{8\sqrt{2}}\left(\sqrt{1+\alpha^2}+\alpha\right)^{\frac{3}{2}}\left\{\left(a_i+b_i\right)^3-3\left(\sqrt{1+\alpha^2}-\alpha\right)^2\left(a_i-b_i\right)^2\left(a_i+b_i\right)\right\}$$

Σ means the summation over one revolution. Putting the angle of rotation of the coordinates as Ψ, given by

$\tan 3\Psi = -B/A$, the coordinates are transformed as

$$\begin{pmatrix} X' \\ Y' \end{pmatrix} = \begin{pmatrix} \cos\Psi & \sin\Psi \\ -\sin\Psi & \cos\Psi \end{pmatrix} \begin{pmatrix} X \\ Y \end{pmatrix}.$$

Then, we can rewrite the Hamiltonian as

$$H = \frac{\varepsilon}{2}\left(X'^2 + Y'^2\right) + \frac{\beta g\sqrt{A^2+B^2}}{3\sqrt{1+\alpha^2}}\left(X'^3 - 3X'Y'^2\right),$$

from which we can obtain following three unstable fixed points:

$$A: (X'_1, Y'_1) = \left(\frac{\sqrt{1+\alpha^2}}{\beta g\sqrt{A^2+B^2}}, 0\right)$$

$$B: (X'_2, Y'_2) = \left(-\frac{1}{2}X'_1, -\frac{\sqrt{3}}{2}X'_1\right)$$

$$C: (X'_3, Y'_3) = \left(-\frac{1}{2}X'_1, \frac{\sqrt{3}}{2}X'_1\right)$$

Transforming back to the (x,x') plane, we obtain the following unstable fixed point:

$$A: (x_1, x'_1) = \left(G\cos(\psi-\chi), -\frac{\sqrt{1+\alpha^2}}{\beta}G\cos(\psi+\chi)\right)$$

$$B: (x_2, x'_2) = \left(G\cos\left(\psi-\chi+\frac{2}{3}\pi\right), -\frac{\sqrt{1+\alpha^2}}{\beta}G\cos\left(\psi+\chi+\frac{2}{3}\pi\right)\right)$$

$$C: (x_3, x'_3) = \left(G\cos\left(\psi-\chi-\frac{2}{3}\pi\right), -\frac{\sqrt{1+\alpha^2}}{\beta}G\cos\left(\psi+\chi-\frac{2}{3}\pi\right)\right)$$

where

$$\tan\chi = \sqrt{1+\alpha^2} + \alpha$$

and

$$G = \frac{\varepsilon\left(1+\alpha^2\right)^{\frac{3}{4}}}{\beta g\sqrt{A^2+B^2}}.$$

The area of the triangular separatrix in the (x,x') plane is given by

$$S = \frac{3\sqrt{3}}{4}\left|\sin 2\chi\right|\frac{\sqrt{1+\alpha^2}}{\beta}G^2.$$

S is assumed to be 10π mm.mrad. The momentum spread is $dp/p = \pm 0.1\%$. So a beam is extracted from that with $dp/p = 0.1\%$ to that with $dp/p = -0.1\%$ by turns. When a beam with $dp/p=0$ has horizontal tune value Q_0, a beam with $dp/p=0.1\%$ has $Q_0+\xi Q_0 dp/p$. If ε is same value, the extracted direction of the beam is same. Result of the simulation is shown in Fig.4. As shown in Fig.4, at the entrance of ES the outgoing trajectories is overlapped each other.

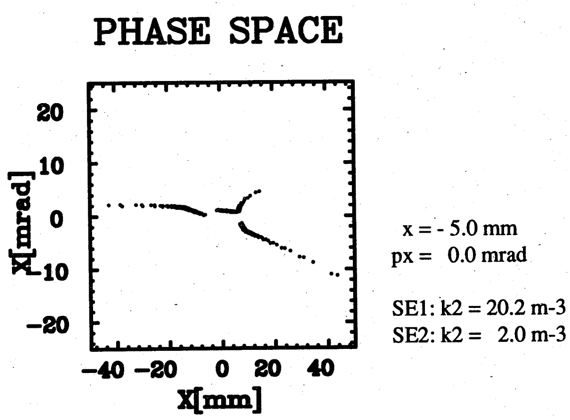

Fig.4 The separatrices and outgoing trajectories of beams with dp/p=-0.1%, 0%, 0.1%, respectively, at ES

4 SUMMARY

We have found that ion beam slow extraction is carried out by using third order resonance (ν_{res}=20/3). This procedure is required to be further investigated in detail for the installation of all the necessary hardwares. Optimization of the lattice of the BSR is under progress.

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